

XXIII Training Course in the Physics of Strongly Correlated Systems

October 4-9, 2021, Salerno

Quantum Simulations using quantum computers on the cloud

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Quantum computing on the cloud



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Research Partners



Financial Support



On sabbatical with



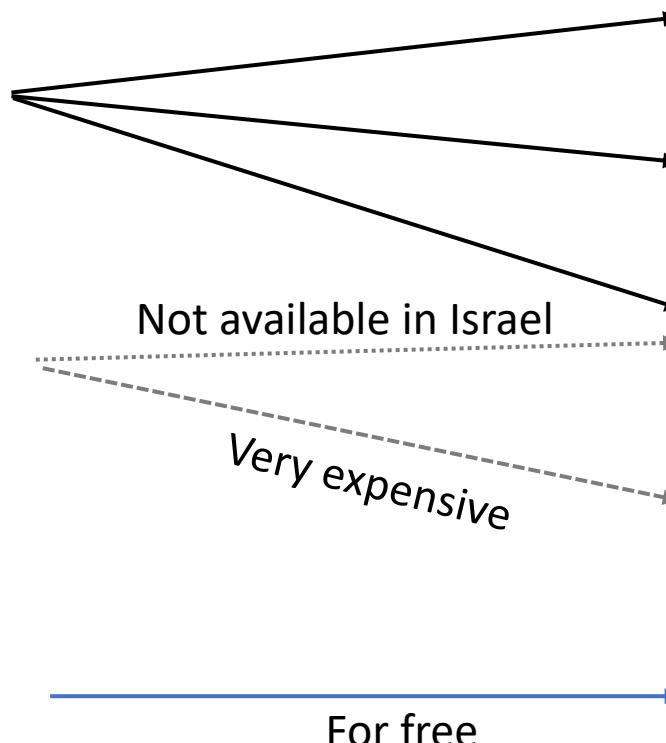
Quantum computing on the cloud is a reality!

Portals:

amazon

Microsoft

IBM



Hardware:

D-Wave

rigetti

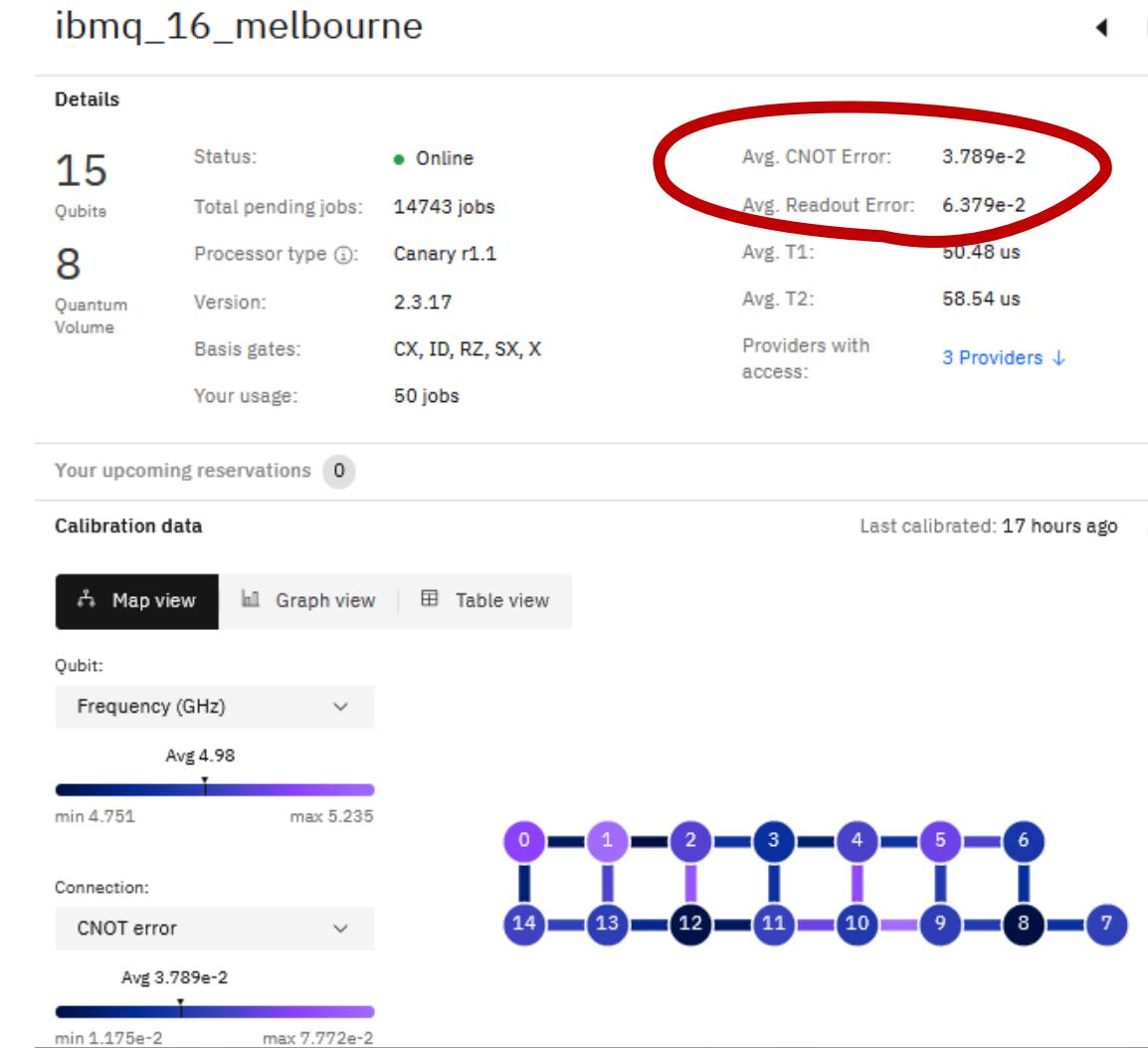
IONQ

Honeywell

IBM



<http://quantum-computing.ibm.com>



Cat state preparation

OpenQASM 2.0 ▾

Open in Quantum Lab

```
1 OPENQASM 2.0;
2 include "qelib1.inc";
3
4 qreg q[3];
5 creg c[3];
6
7
```

H \oplus \otimes \oplus \ominus \otimes \ominus I T S Z T^\dagger S^\dagger P RZ \bullet $|0\rangle$ \otimes^z ⓘ : OpenQASM 2.0 ▾

if | \sqrt{X} \sqrt{X}^\dagger Y RX RY U RXX RZZ + Add

q₀

q₁

q₂

+

Please wait....

Probabilities ▾

Probability (%)

Computational basis states	Probability (%)
000	100
001	0
010	0
011	0
100	0
101	0
110	0
111	0

Statevector ▾

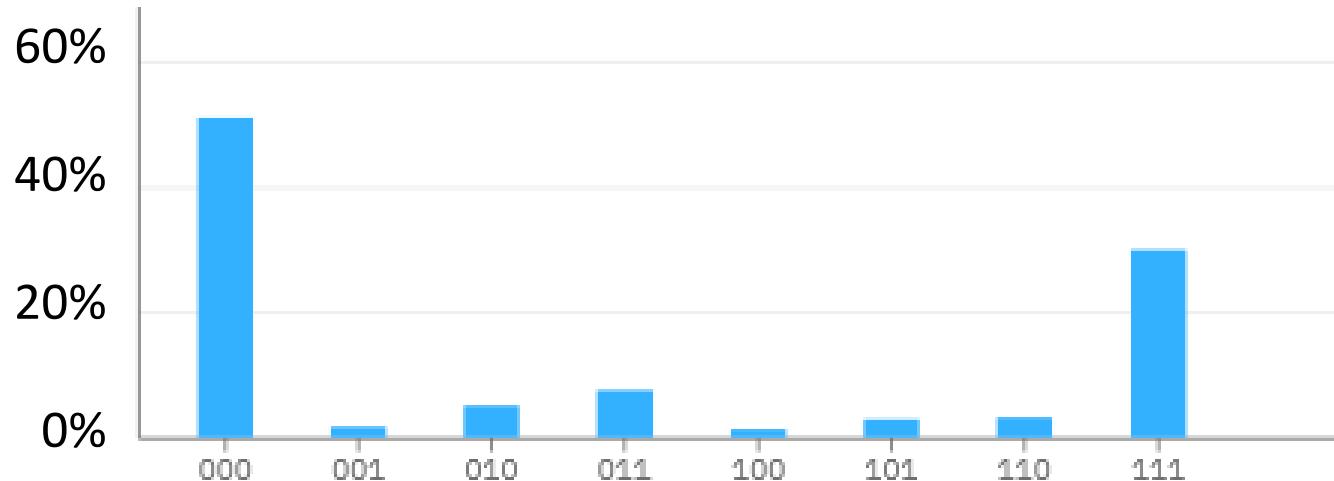
Amplitude

Computational basis states	Amplitude
000	1
001	0
010	0
011	0
100	0
101	0
110	0
111	0

E I



Cat state preparation



Demonstration of Shor's factoring algorithm for $N=21$ on IBM quantum processors

Unathi Skosana and Mark Tame*

Department of Physics, Stellenbosch University, Matieland 7602, South Africa

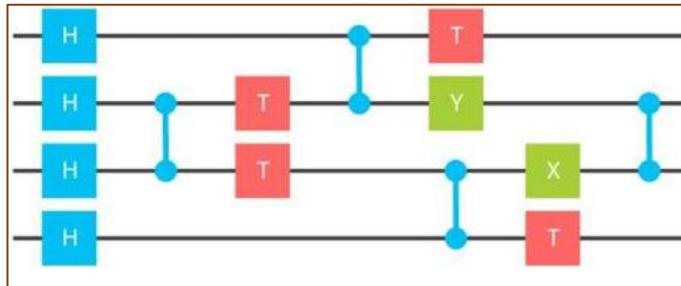
(Dated: March 26, 2021)

We report a proof-of-concept demonstration of a quantum order-finding algorithm for factoring the integer 21. Our demonstration involves the use of a compiled version of the quantum phase estimation routine, and builds upon a previous demonstration by Martín-López et al. in *Nature Photonics* 6, 773 (2012). We go beyond this work by using a configuration of approximate Toffoli gates with residual phase shifts, which preserves the functional correctness and allows us to achieve a complete factoring of $N = 21$. We implemented the algorithm on IBM quantum processors using only 5 qubits and successfully verified the presence of entanglement between the control and work register qubits, which is a necessary condition for the algorithm's speedup in general. The techniques we employ may be useful in carrying out Shor's algorithm for larger integers, or other algorithms in systems with a limited number of noisy qubits.

The main challenge

Model :

Unitary quantum computer



f.e. Shor algorithm (breaks RSA),
quantum machine learning

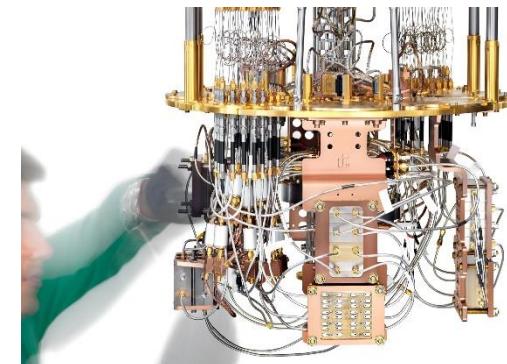
Reality :

Noisy superconducting circuits

quantum error correction

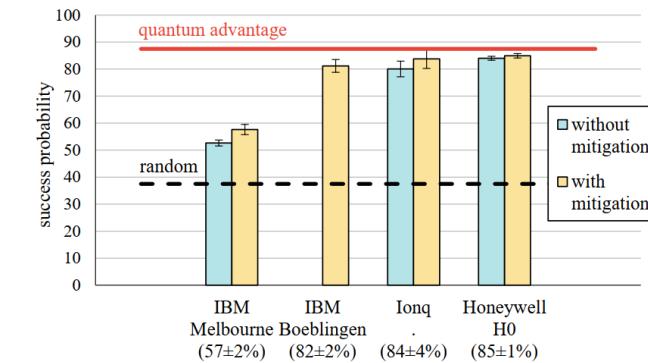


open quantum systems

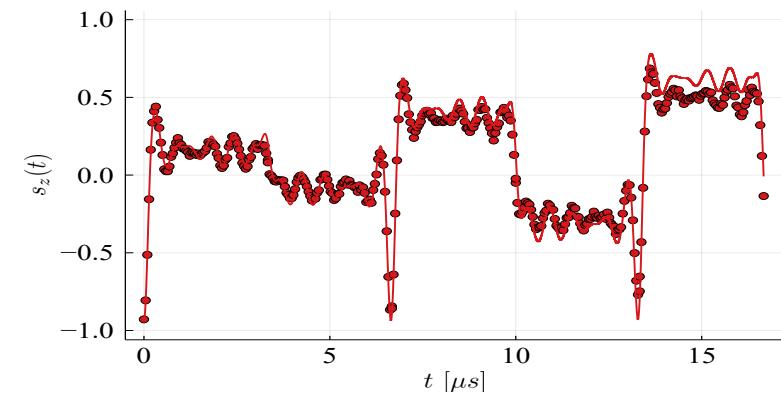


Three examples of quantum simulations

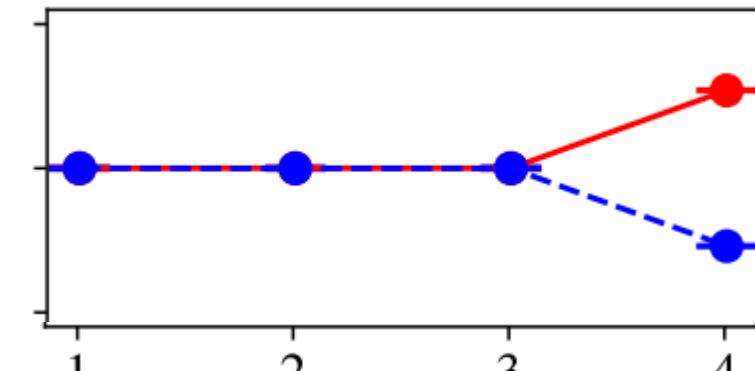
1. Noisy nonlocal games arXiv:2105.05266



2. Floquet engineering arXiv:2102.09590



3. Topological phases arXiv:2002.04620 (PRL, 2020) arXiv:2008.09332 (PRB, 2020)



Bell inequality... with many qubits

a.k.a. Non-local quantum games

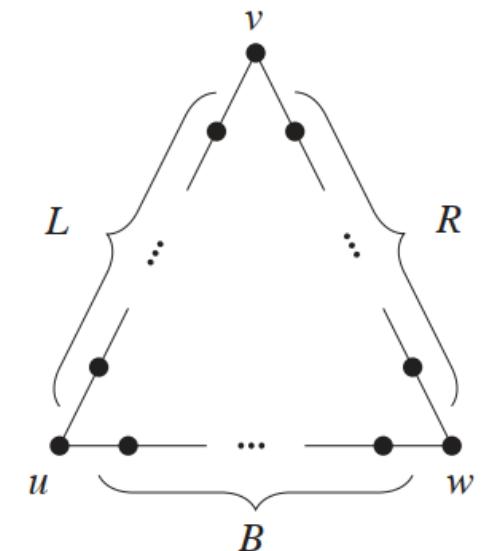
Example: triangle game using a 1d cluster state

$$\langle P(\text{win}) \rangle_{\text{quantum}} = 1$$

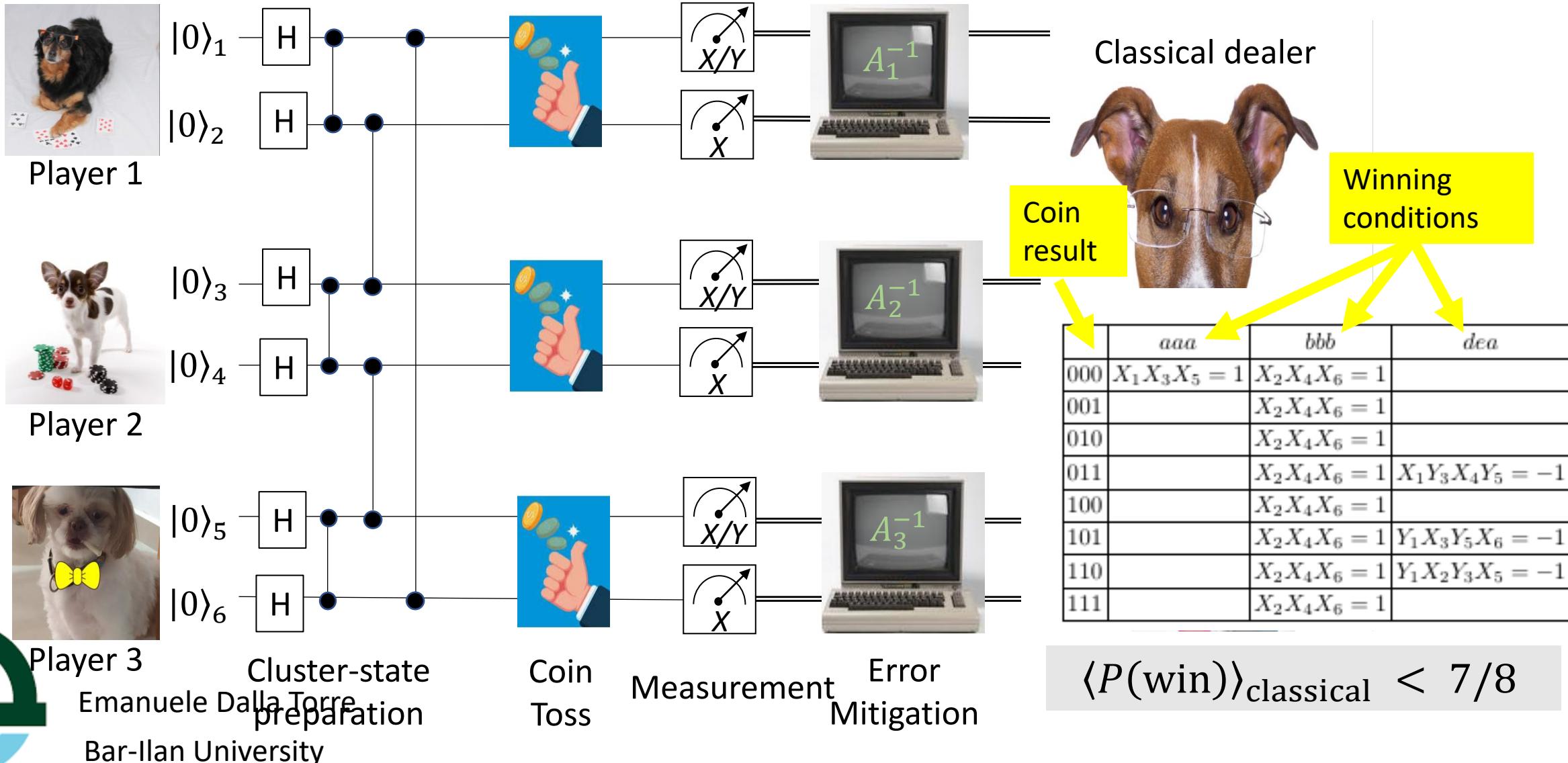
$$\langle P(\text{win}) \rangle_{\text{classical}} < 7/8$$

Bravyi, Gosset & Konig, Science 2018

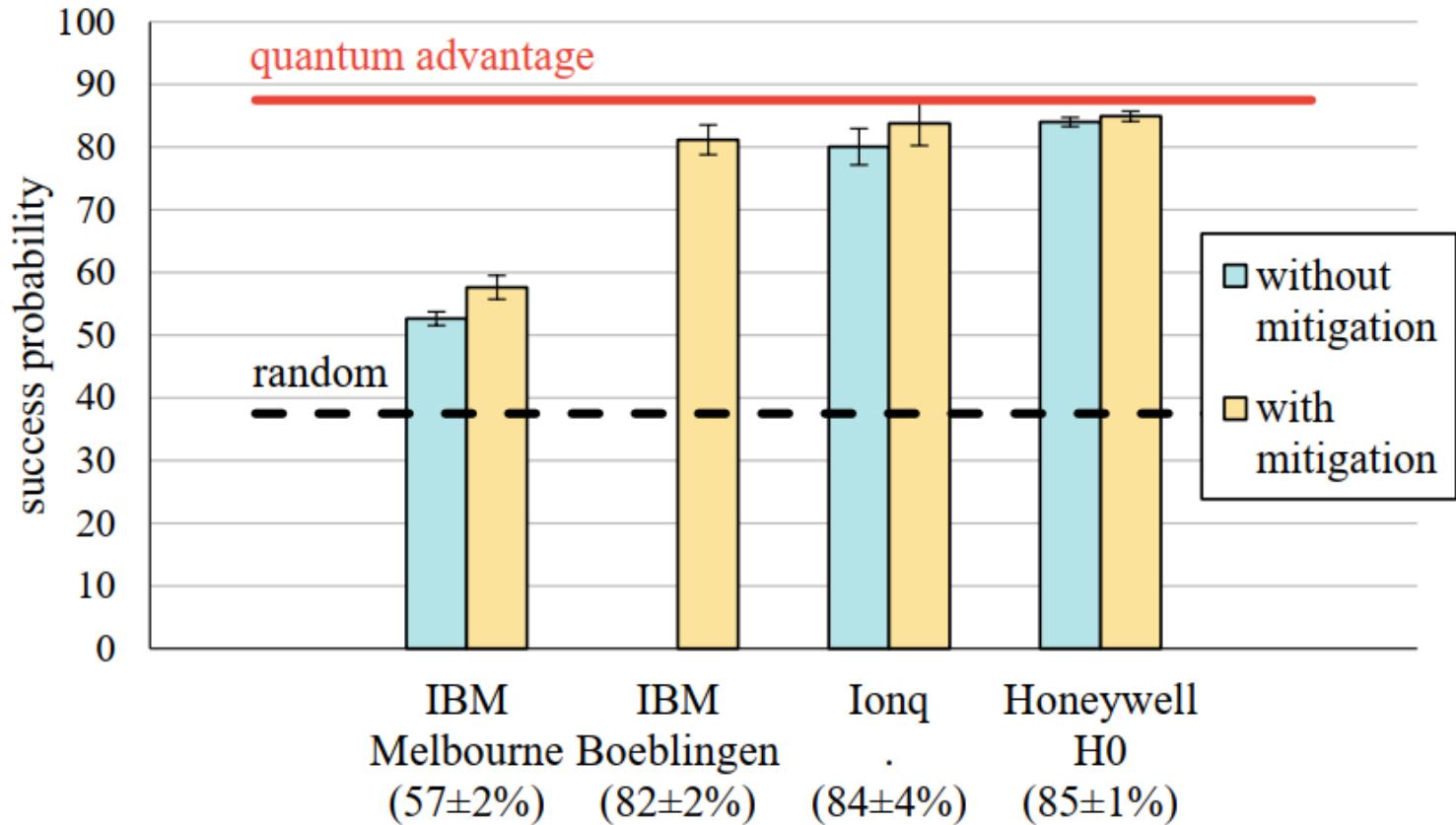
Daniel & Myiaka, PRL 2021



Minimal realization : 3 players = 6 qubits



Triangle game (6 qubit) : results



If all the students fail a test... lower the bar!

Stabilizers

$$s_i = X_{i-1} Z_i X_{i+1}$$

$$s_i |\psi_{\text{cluster}}\rangle = |\psi_{\text{cluster}}\rangle$$

Sum of all products

$$S_{\text{all}} = 1 + \sum_i s_i + \sum_{i,j} s_i s_j + \dots$$

$$S_{\text{all}} |\psi_{\text{cluster}}\rangle = 2^n |\psi_{\text{cluster}}\rangle$$

$$(S_{\text{all}})_{\text{classic}, n=6} \leq 28$$

Guhne, Toth, Hyllus, Briegel PRL (2005)

Optimal sum

$$S_{\text{optimal}} = \sum_{i,j} s_i s_j + \sum_{i,j,k} s_i s_j s_k + \sum_{i,j,k,l} s_i s_j s_k s_{k+1}$$

$$\langle S_{\text{optimal}} \rangle_{\text{IonQ}} = 41 \pm 0.5$$

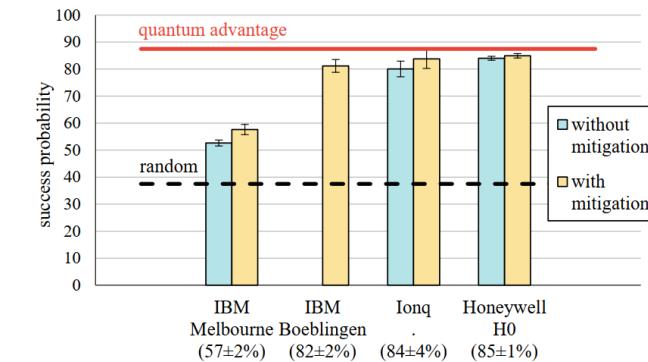
$$(S_{\text{optimal}})_{\text{classic}, n=6} \leq 19$$

Cabello, Guhne, Rodriguez PRA (2008)

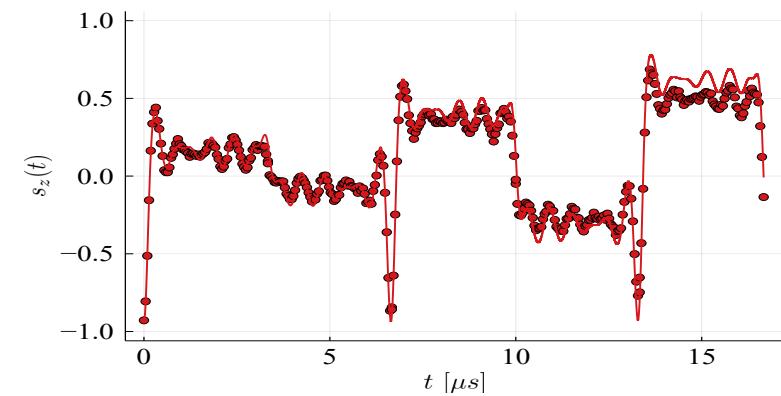


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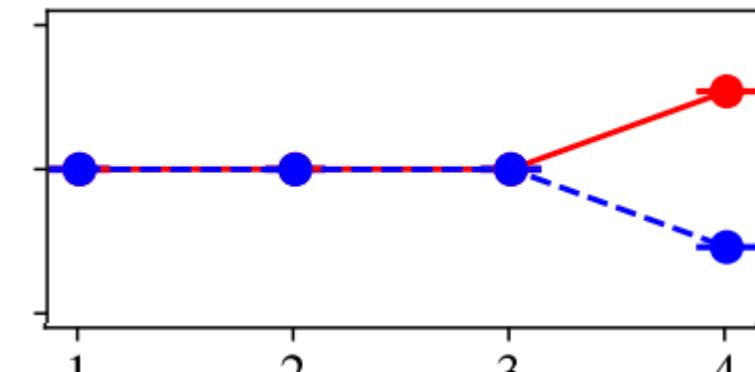
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Periodically driven single qubit

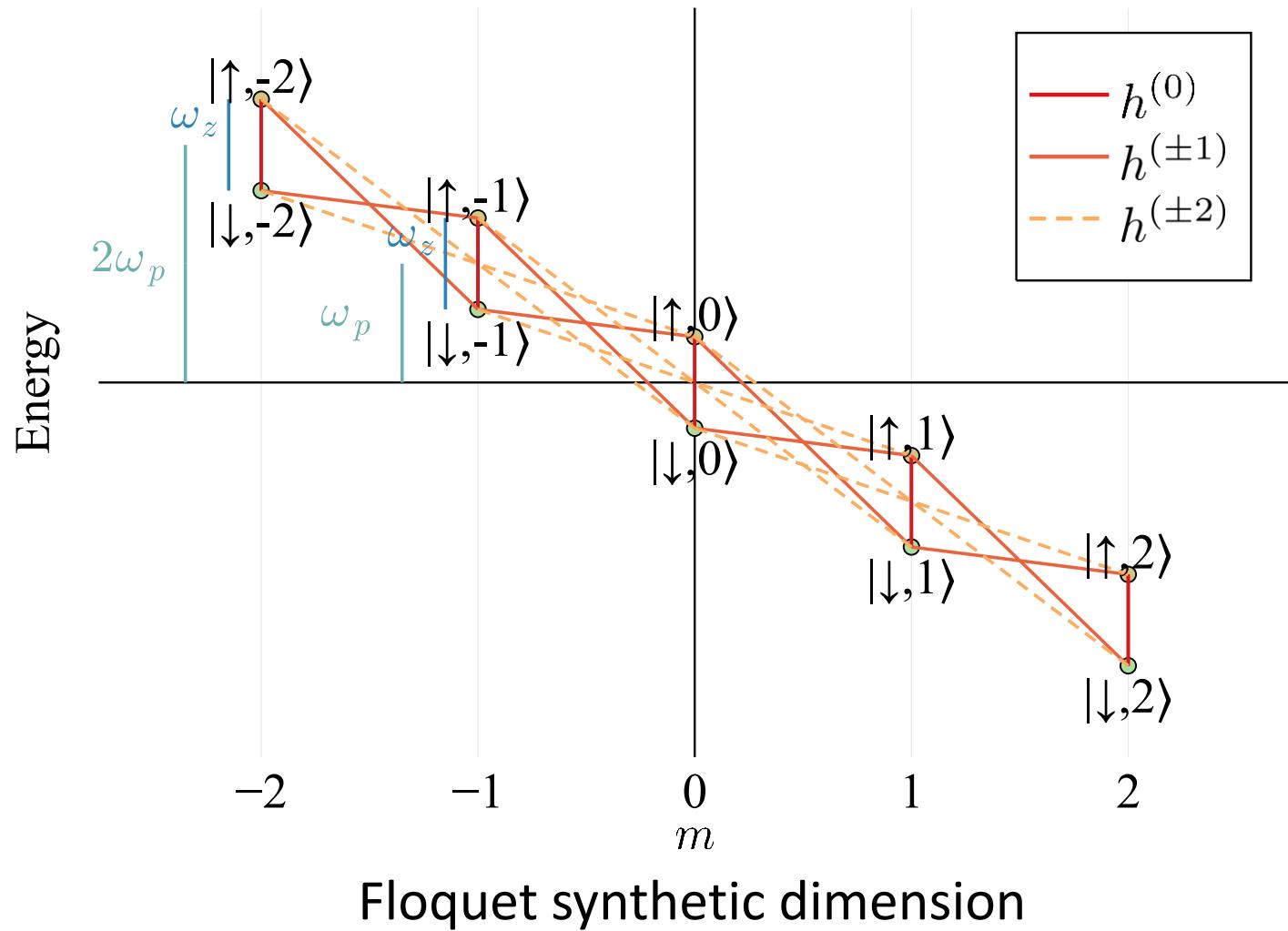
Main ingredient:

Pulse engineering:

$$H(t) = \omega_z \sigma^z + h(t) \sigma^x$$

Periodic drive (Floquet)

$$h(t + \tau) = h(t) \quad \omega_p = \frac{2\pi}{\tau}$$

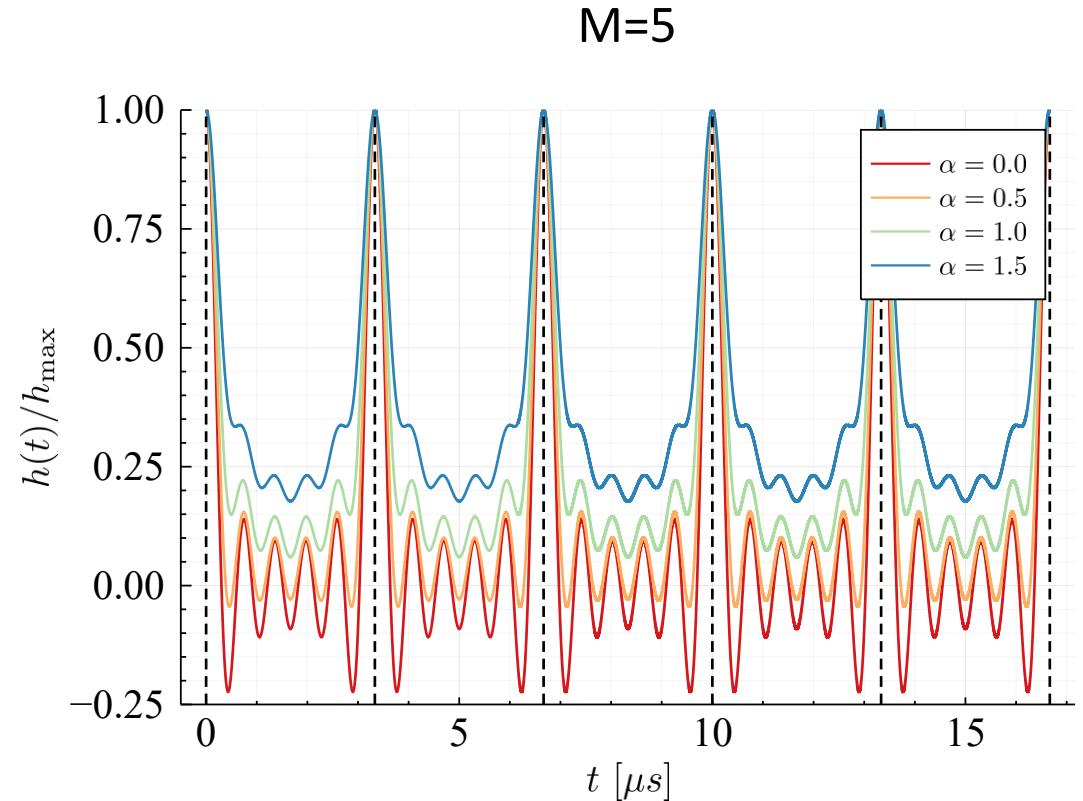


Floquet engineering

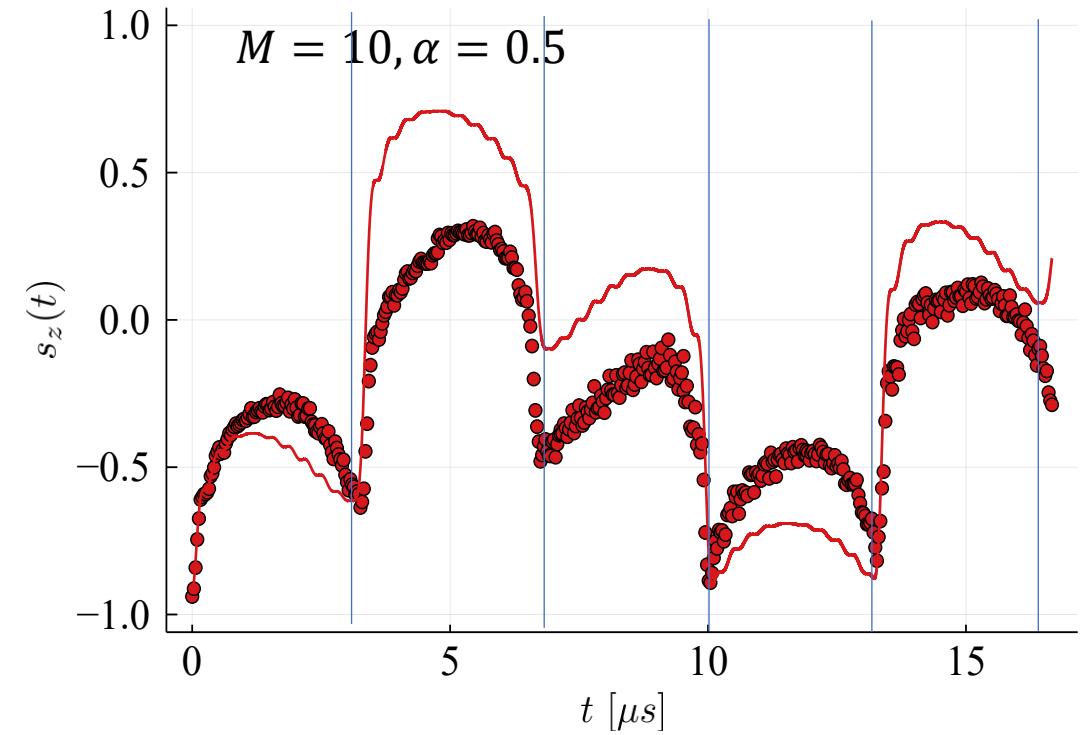
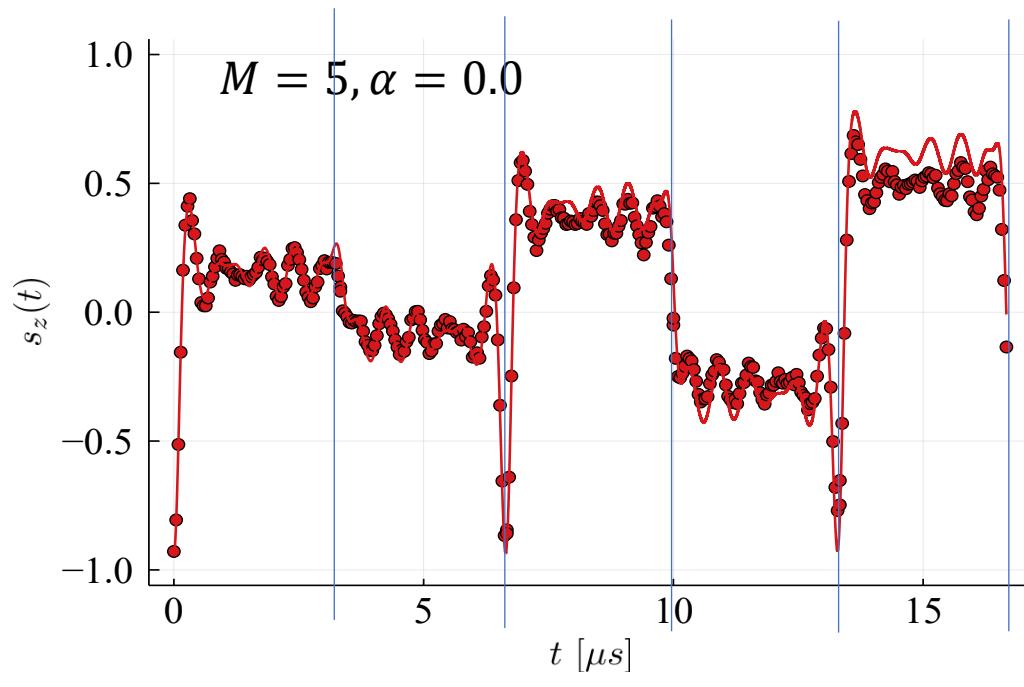
Long ranged coupling:

$$h^{(m)} = \frac{h_0}{(1 + |m|)^\alpha}$$

Simulation of gravitational models &
unscreened Coloumb forces ($\alpha = 1$)!

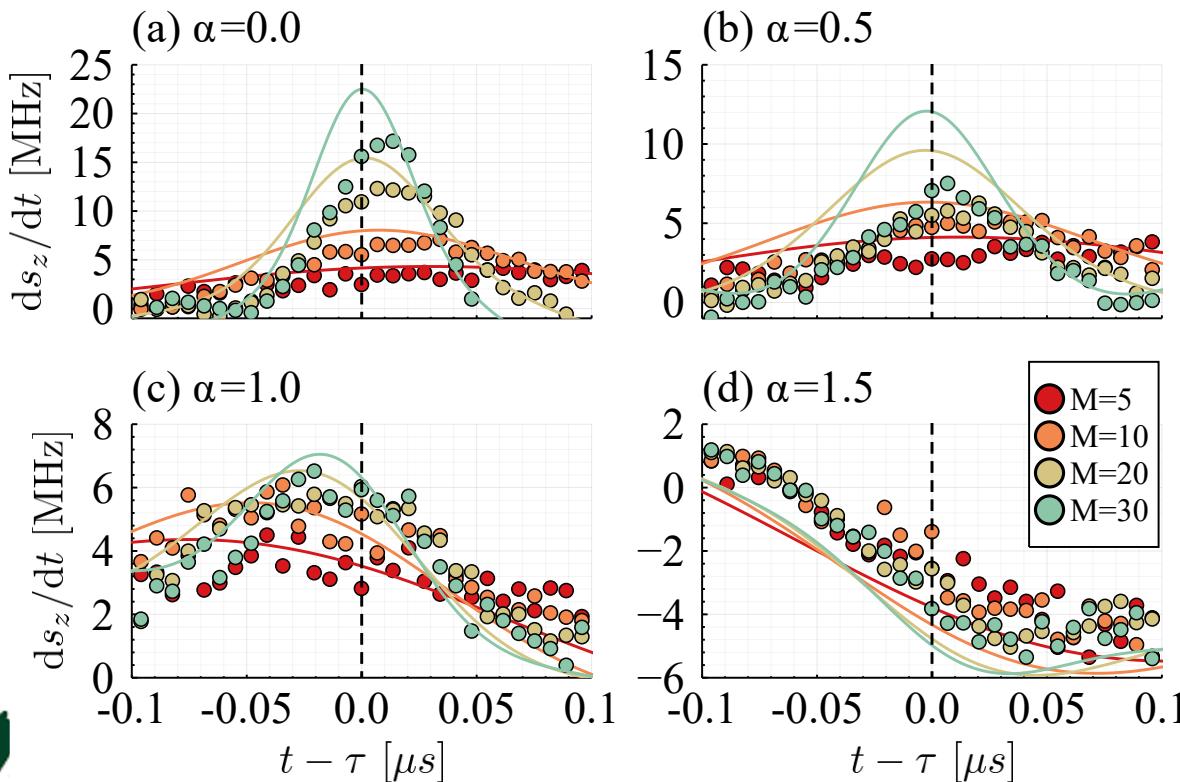


Result from a physical device

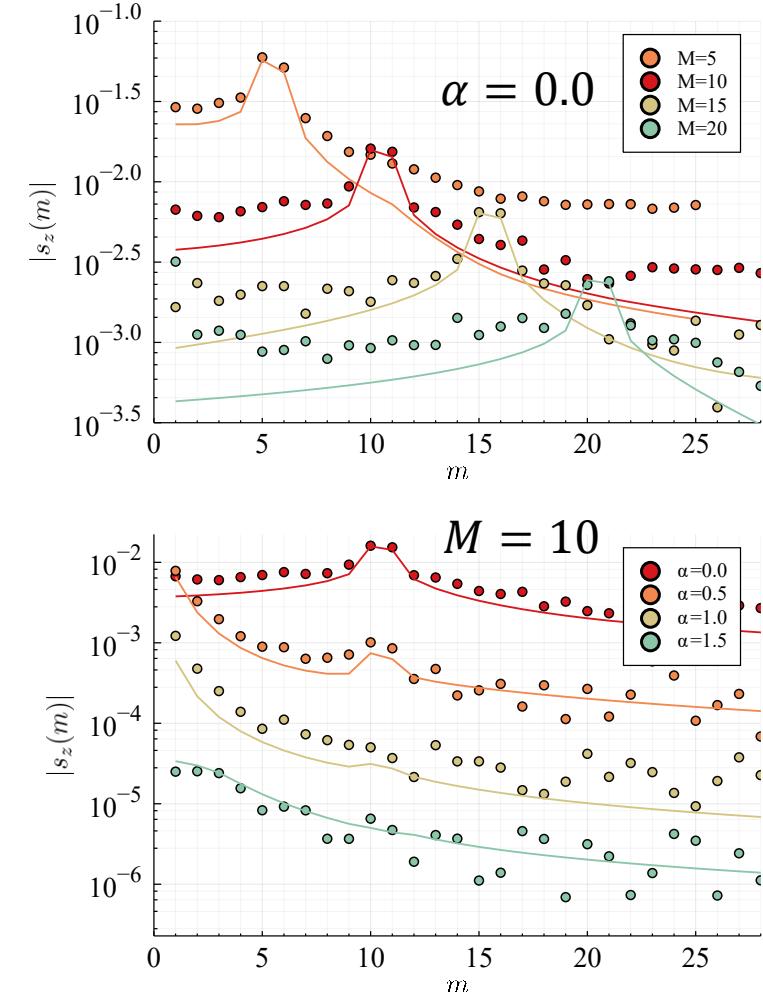


Transition from short range to long range

Time derivatives: $\frac{ds_z}{dt}(t = \tau) \sim \sum_m \frac{1}{|m|^\alpha}$

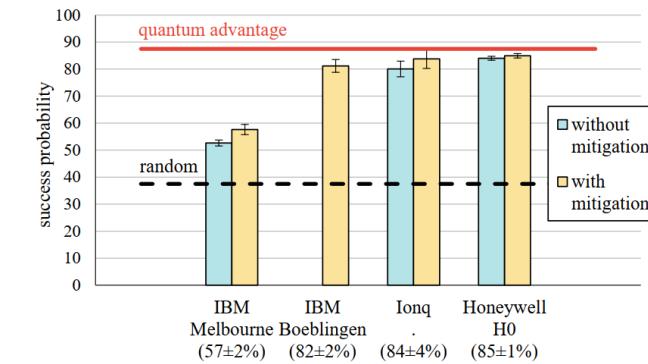


Long tails of the eigenstates:

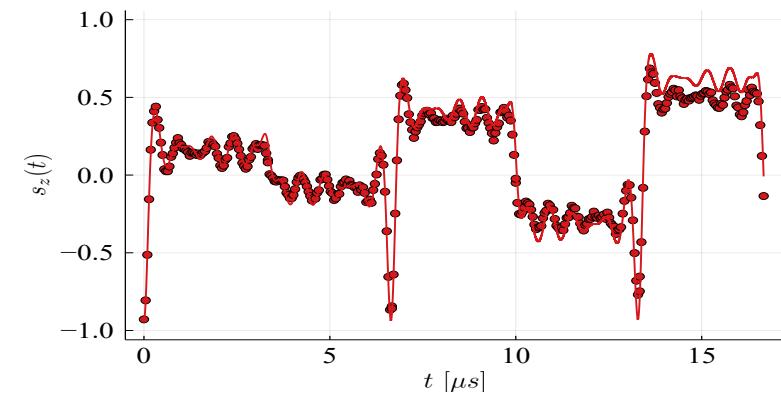


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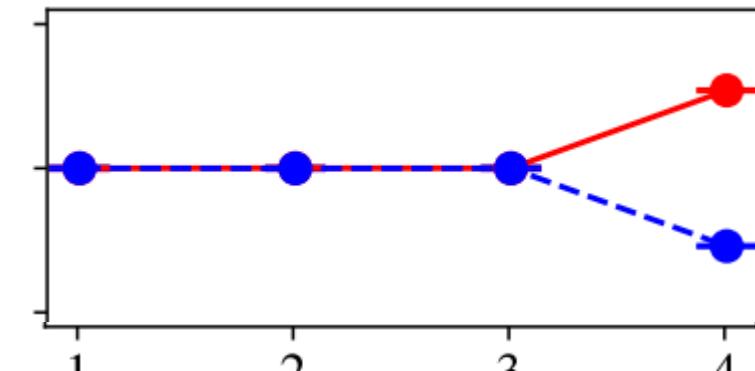
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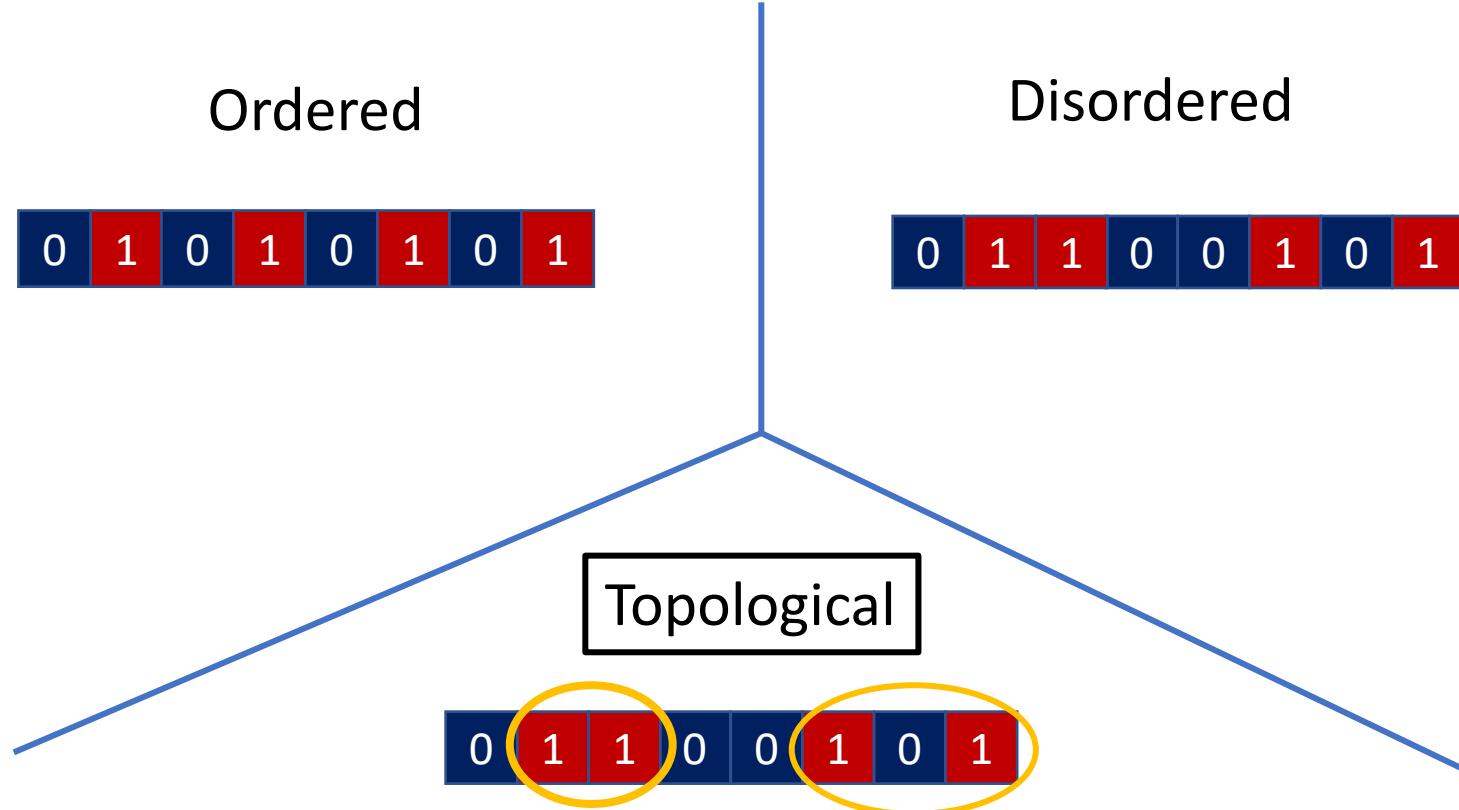
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Quantum phases of matter

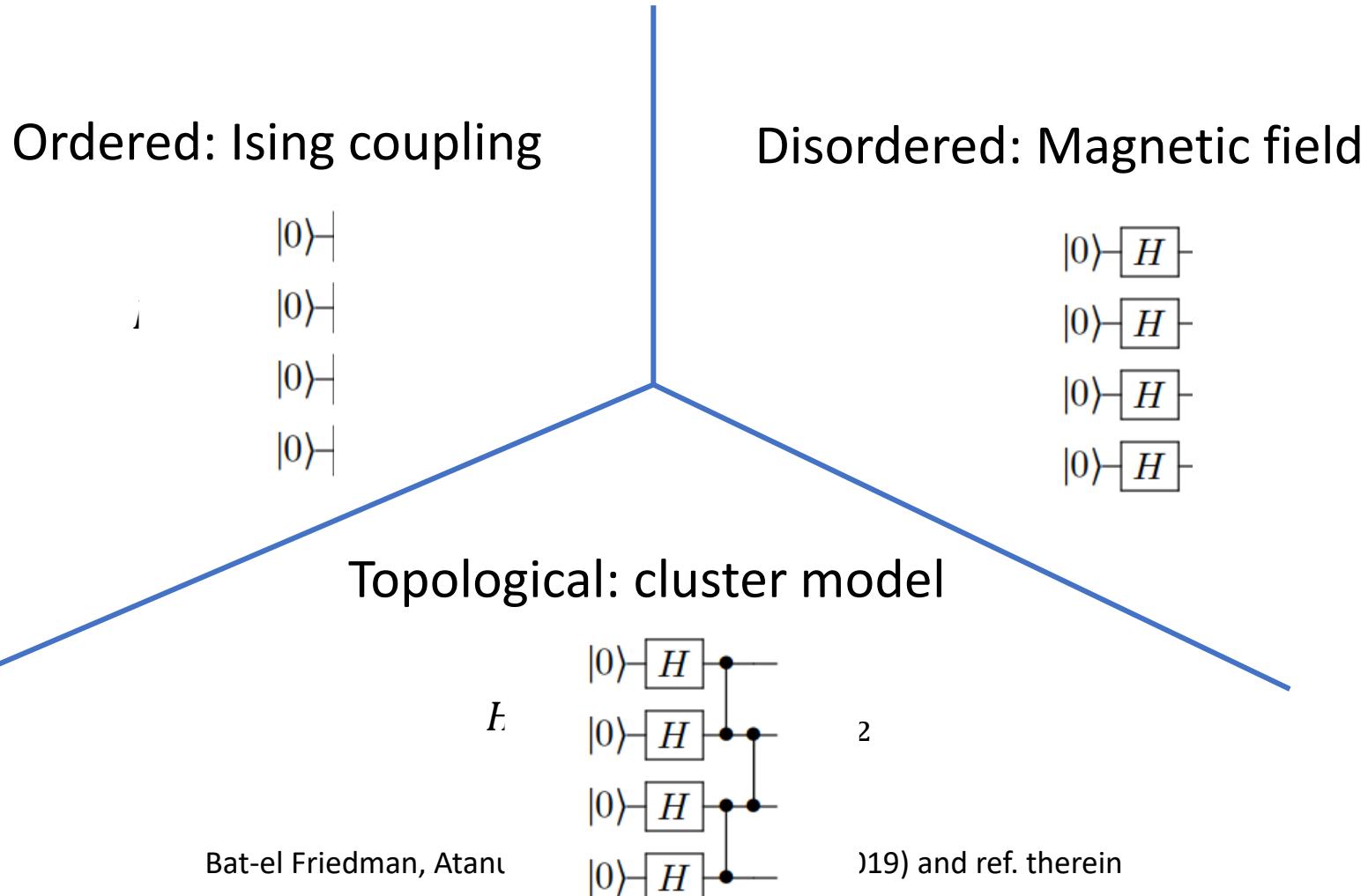


Haldane (1982), AKLT (1987), den Nijs&Rommelse (1989), Dalla Torre,Berg&Altman (2006),



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Quantum phases of matter



Symmetry resolved topological (SPT) phases

Pollmann-Turner-Berg-Oshikawa (2011), Chen-Gu-Liu-Wen (2012)

$$\text{Symmetry} \rightarrow G |\psi\rangle = g|\psi\rangle \rightarrow [G, \rho_A] = 0$$

Symmetry resolved
reduced density matrix:

$$\rho_A = \begin{cases} \text{Trivial} & \tilde{\rho}_A(+) \\ & 0 \\ & 0 \quad \tilde{\rho}_A(-) \end{cases}$$
$$\rho_A = \begin{cases} \text{SPT} & \tilde{\rho}_A(+) \quad 0 \\ & 0 \quad \tilde{\rho}_A(-) \end{cases}$$

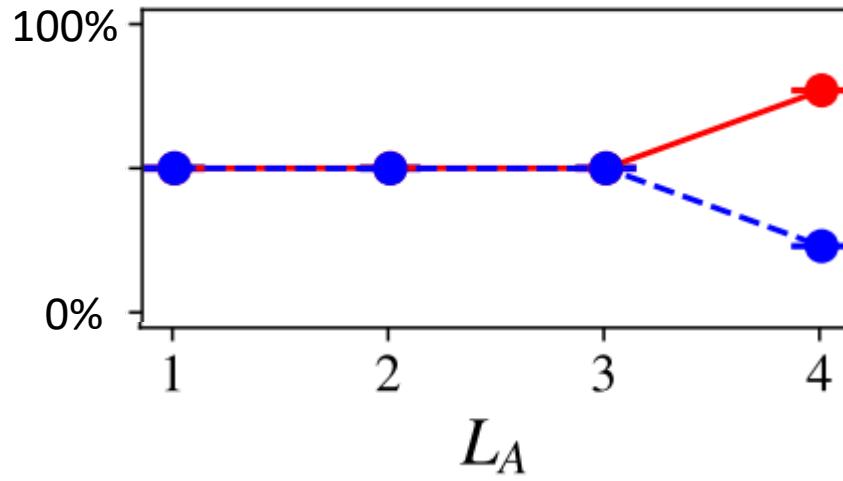
Goldstein & Sela (PRL, 2018)



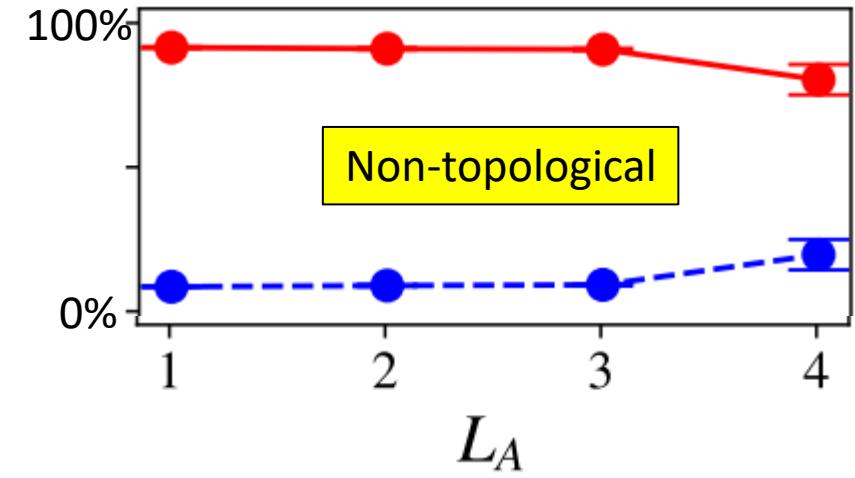
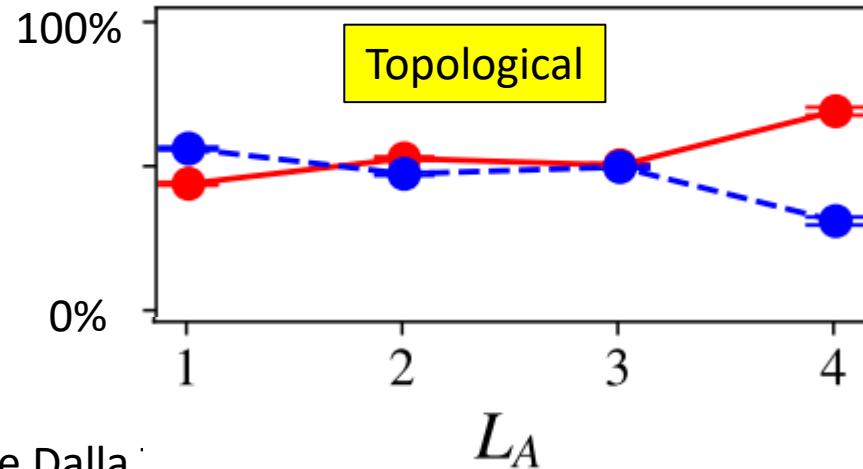
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Identifying a topological phases

Simulation



Experiment

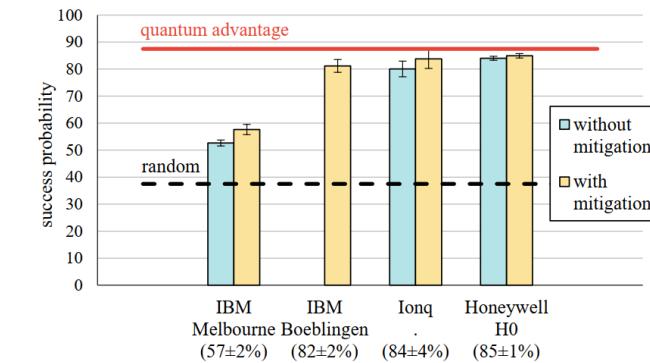


Emanuele Dalla

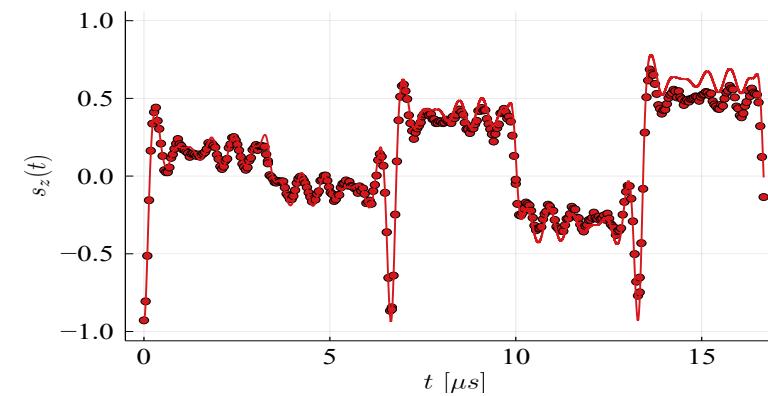
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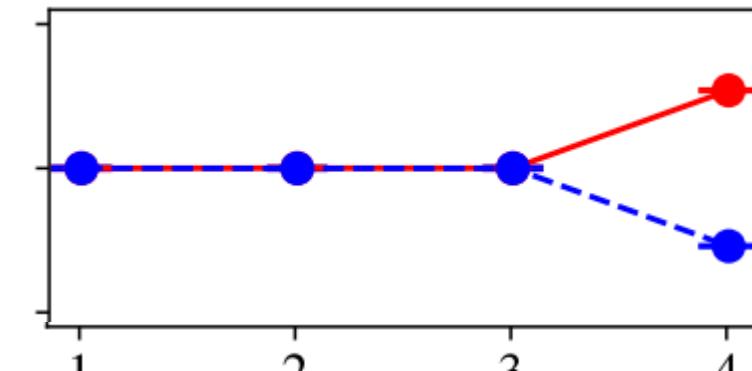
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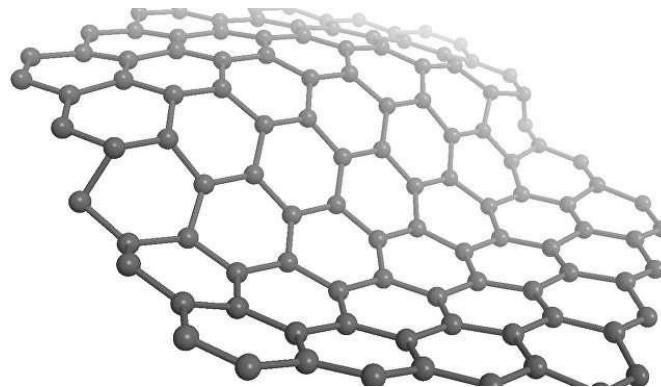
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Summary: Quantum simulations

Quantum molecules/material

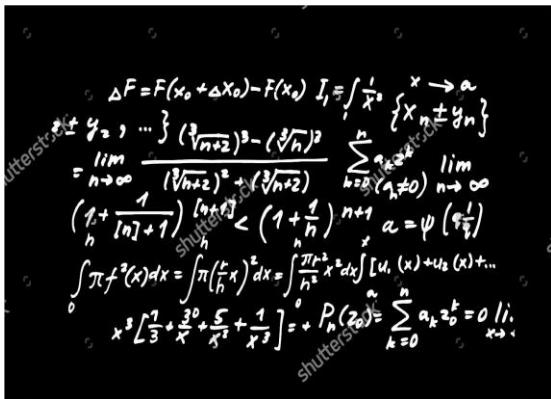
Size of Hilbert space = 2^{N_A}



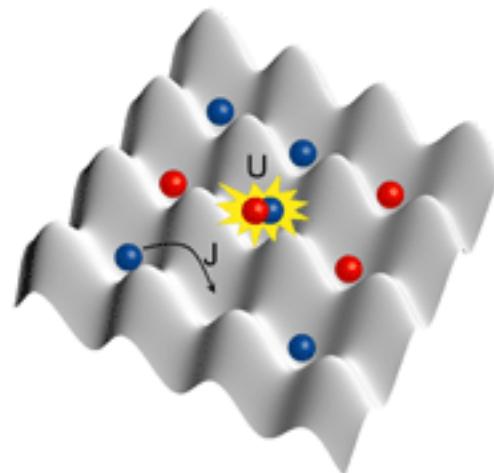
Classical supercomputers



Quantum field theories


$$\Delta F = F(x_0 + \Delta x_0) - F(x_0)$$
$$I_1 = \int \frac{1}{x^2} \{x \rightarrow a, x \pm y_n\}$$
$$= \lim_{n \rightarrow \infty} \frac{(3\sqrt{n+2})^3 - (3\sqrt{n})^3}{(3\sqrt{n+2})^3 + (3\sqrt{n+2})}$$
$$\sum_{n=0}^{\infty} \lim_{\substack{x \rightarrow a \\ b \rightarrow 0}} \lim_{n \rightarrow \infty}$$
$$\left(\frac{1}{n} + \frac{1}{[n] + 1}\right)^{3/2} < \left(1 + \frac{1}{n}\right)^{n+1} a = \psi\left(\frac{1}{n}\right)$$
$$\int \pi f^2(x) dx = \int \pi \left(\frac{r}{h} x\right)^2 dx = \int \frac{\pi r^2}{h^2} x^2 dx [u_r(x) + u_s(x) + \dots]$$
$$x^2 \left[\frac{r^2}{3} + \frac{3r}{h} + \frac{5}{h^2} + \frac{1}{x^2}\right] = + P_n(x) = \sum_{k=0}^n a_k x^k = 0 / i.$$

Ultracold atoms

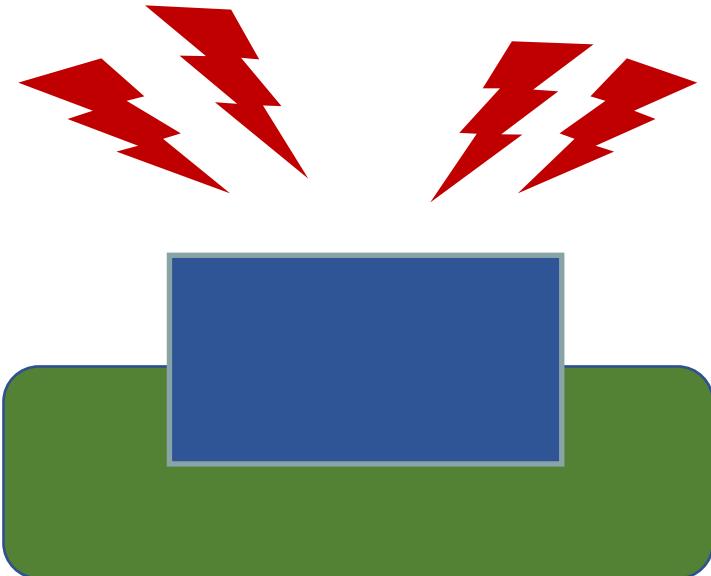


Quantum computer on the cloud



Quantum simulations with quantum computers on the cloud

Time dependent drive



Quantum system

Environment

Floquet engineering
arXiv:2102.09590

Topological phases
arXiv:2002.04620 (PRL, 2020)
arXiv:2008.09332 (PRB, 2020)

Noisy nonlocal games and
error mitigation
arXiv:2105.05266



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EXTRA SLIDES



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