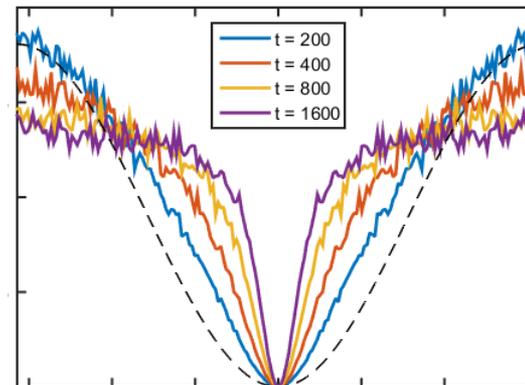
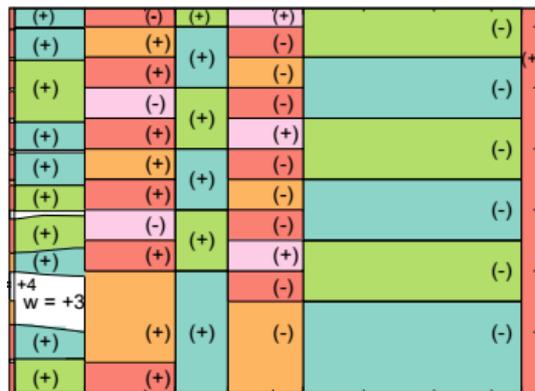
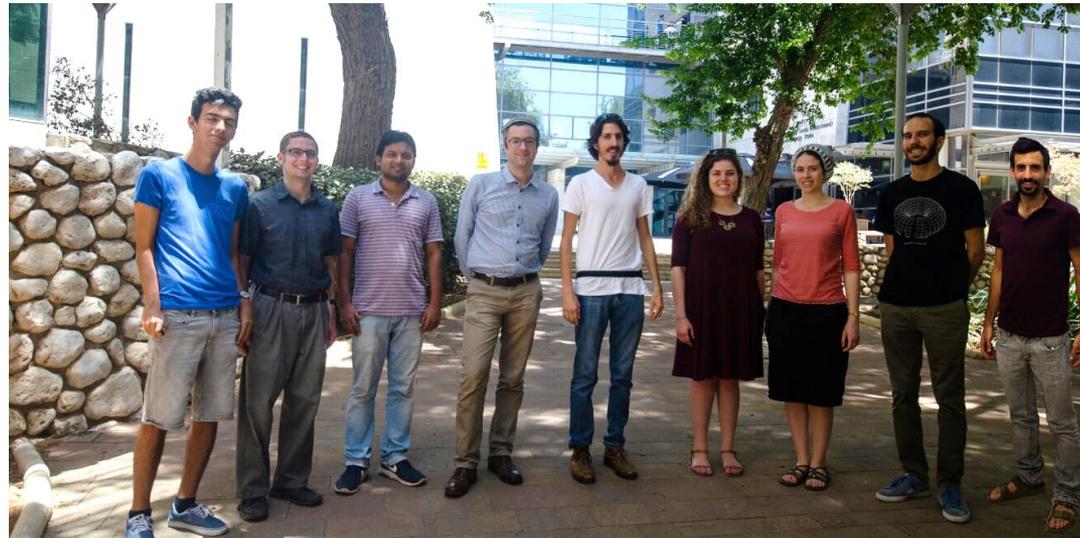


Emanuele Dalla Torre

From Floquet Engineering to Pre-thermalization



Many-body quantum dynamics group



Angelo Russomanno
(→ ICTP, Trieste)



Atanu Rajak
(→ Kolkata)



Itzhack Dana
Bar-Ilan University

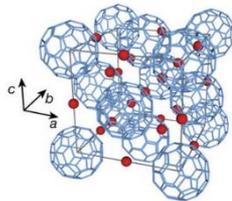


Roberta Citro
University of Salerno

Why Floquet Engineering ?

Superconductors

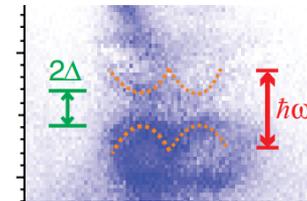
→ Light-enhanced superconductivity



[Cavalleri group – Nature 2016]

Topological insulators

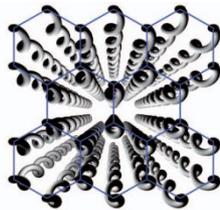
→ Floquet band structure



[Gedik group – Science 2013]

Wave guides

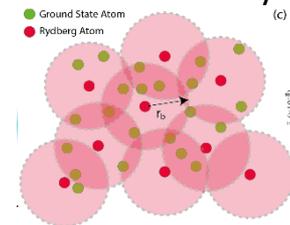
→ Photonic topological insulators



[Segev-Szameit group – Science 2018]

Rydberg atoms

→ Discrete time crystals



[Lukin group – Nature 2017]

Floquet Engineering - idea

Periodically driven Hamiltonian

$$H(t) = H(t + T)$$



Stroboscopic time evolution

$$|\psi(nT)\rangle = (U(T))^n |\psi_0\rangle$$



Effective Floquet Hamiltonian

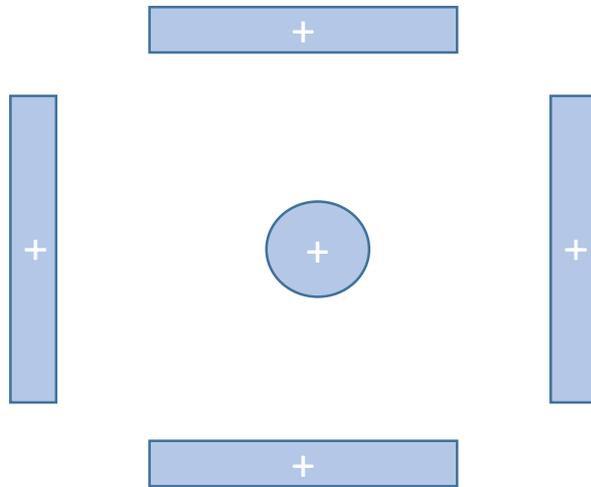
$$U(T) = e^{-i H_F T}$$

$$|\psi(nT)\rangle = e^{-i H_F n T} |\psi_0\rangle$$

H_F is cooooool!

Eckardt, RMP 2017

Example 1 – Paul trap



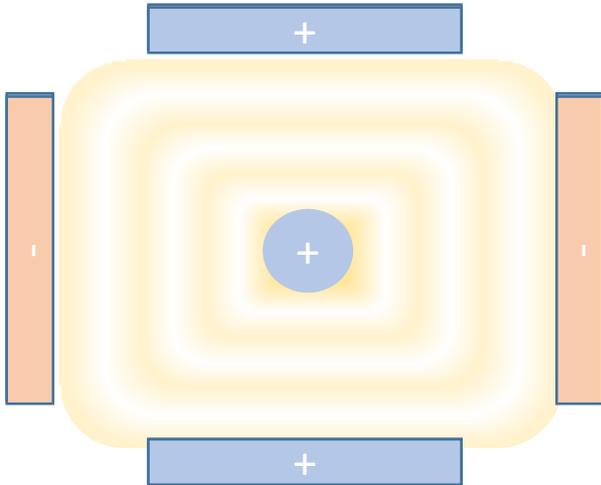
$$\text{Trapping : } \begin{cases} \vec{\nabla} \phi = 0 \\ \nabla^2 \phi > 0 \end{cases}$$

But...

$$\nabla^2 \phi = -\vec{\nabla} \cdot \vec{E} = -4\pi\rho = 0$$

→ No static trapping in vacuum

Example 1 – Paul trap



$$H(t) = \frac{p^2}{2m} + V(x^2 - y^2) \cos(\Omega t)$$

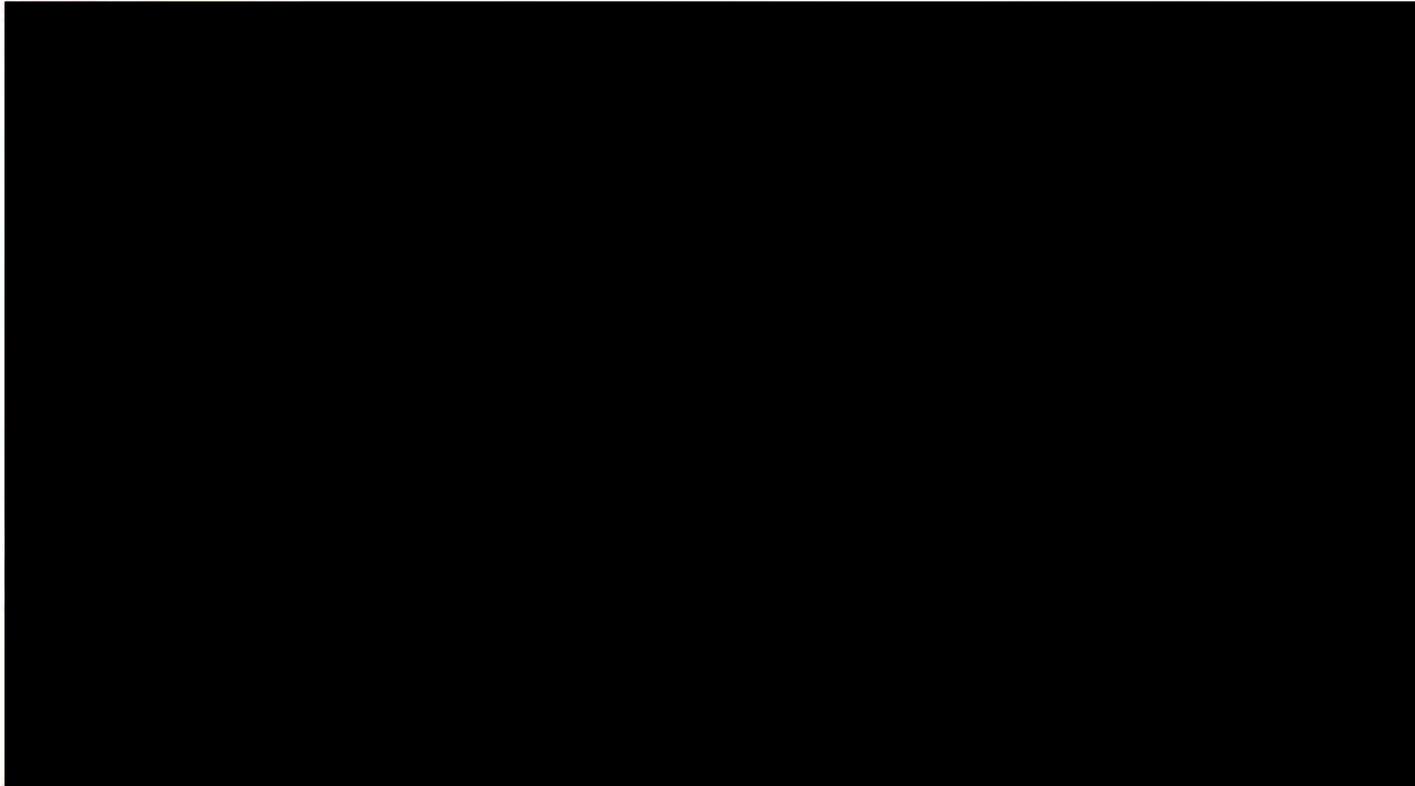
$$H_{av} = \frac{p^2}{2m}$$

Magnus expansion

$$H_F = H_{av} + \frac{1}{\Omega} \int dt \int dt' [H(t), H(t')] + \dots$$

$$H_F = \frac{p^2}{2m} + \frac{V^2}{\Omega^2 m} (x^2 + y^2)$$

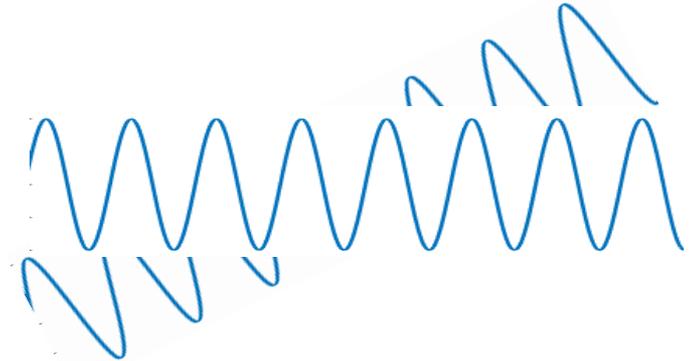
Example 1 – Paul trap



Video credit: [Harvard Natural Sciences Lecture Demonstrations](#)

Example 2 – lattice shaking

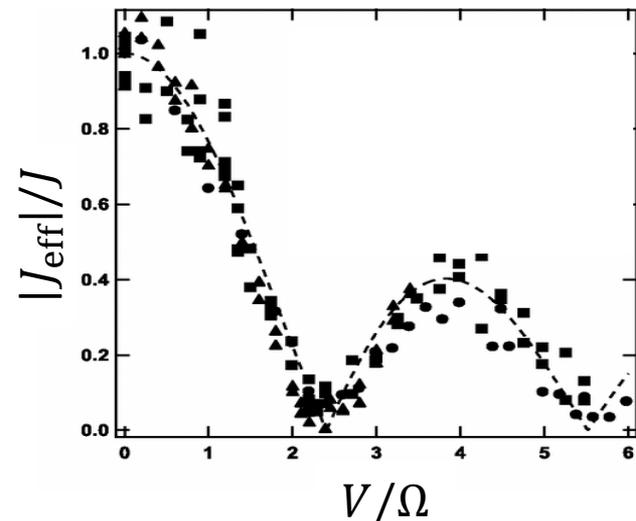
$$H(t) = \sum_i J(c_i^\dagger c_{i+1} + H.c.) + V \sum_i n_i \cos(\Omega t)$$



“Gauge transformation” (rotating frame)

$$H(t) = J \sum_i (e^{iV \cos(\Omega t)} c_i^\dagger c_{i+1} + H.c.)$$

$$H_F \approx H_{av} = \underbrace{J J_0 \left(\frac{V}{\Omega} \right)}_{J_{\text{eff}}} \sum_i (c_i^\dagger c_{i+1} + H.c.)$$



Morsch-Arimondo PRL 2007 & PRL 2008

Many-body periodically driven systems

- **Periodically driven Ising model** → **Time crystal**

Russomanno, Dalla Torre (EPL 2016)

Russomanno, Friedman, Dalla Torre (PRB 2017)

- **Many body kicked rotor** → **Prethermalization**

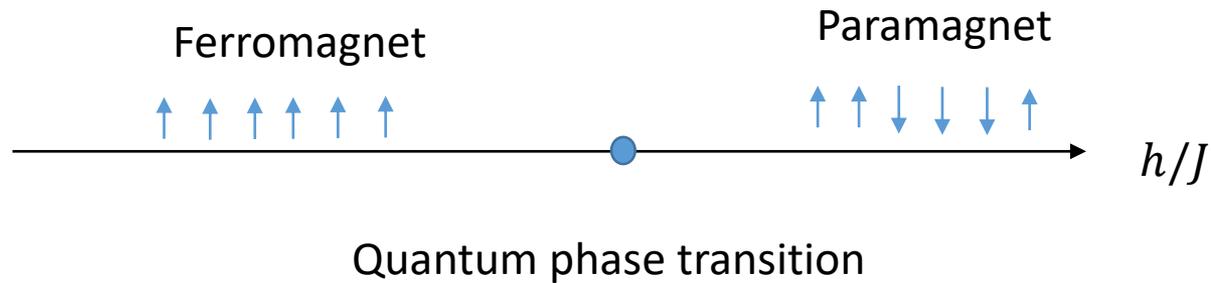
Citro, Dalla Torre, et al, Annals of Physics (2015)

Rajak, Citro, Dalla Torre (J. Phys. A: Math. Theor. 51 465001 (2018))

Rajak, Dana, Dalla Torre (in prep.)

Ising model : static

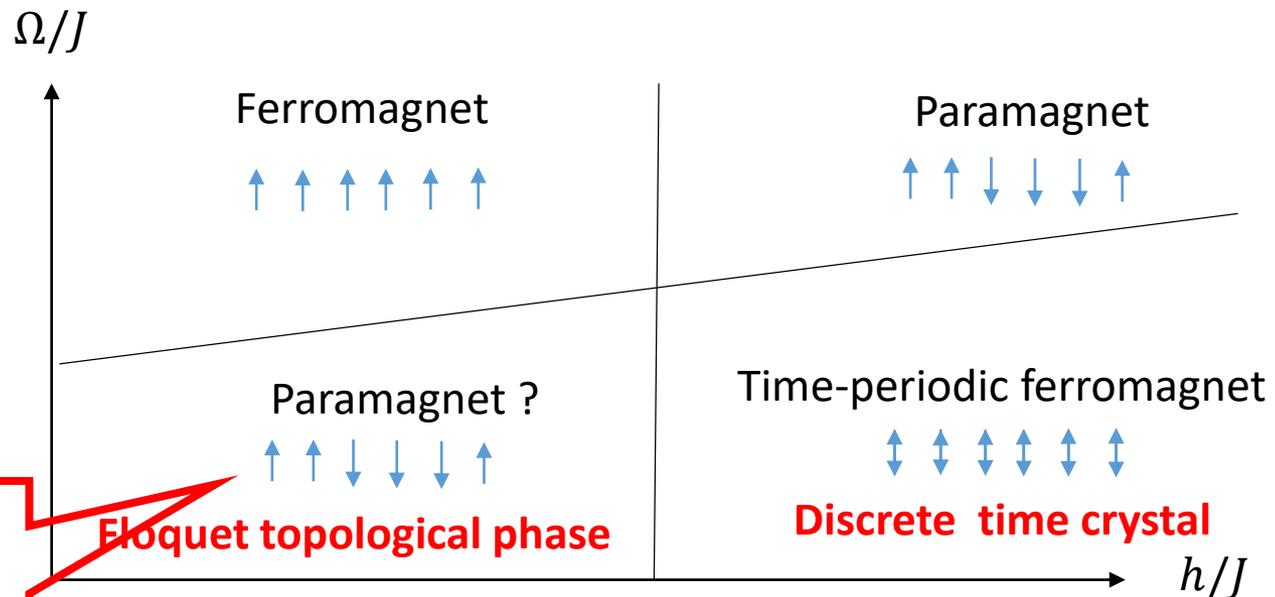
$$\hat{H}(t) = \sum_j (-J\hat{\sigma}_j^z \hat{\sigma}_{j+1}^z + h(t)\hat{\sigma}_j^x) \quad h(t) = h_0$$



Ising model : periodic drive

$$\hat{H}(t) = \sum_j (-J \hat{\sigma}_j^z \hat{\sigma}_{j+1}^z + h(t) \hat{\sigma}_j^x)$$

$$h(t) = h_0 + \delta h \cos(\Omega t)$$



Non local order

$$\hat{\beta}_j^y = \left(\prod_{k=1}^{j-1} \hat{\sigma}_k^x \right) \hat{\sigma}_j^z \hat{\sigma}_{j+1}^y$$

Floquet topological phase

Discrete time crystal

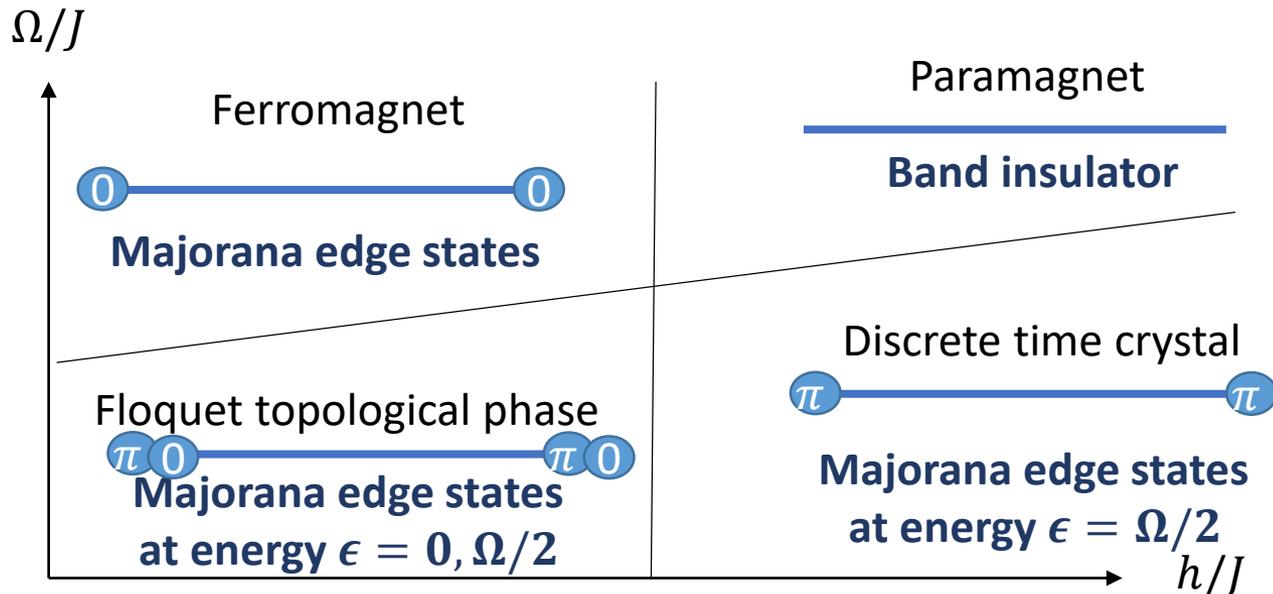
Else, Bauer, Nayak (PRL 2016)

Russomanno, Friedman, Dalla Torre (PRB 2017)

Ising model : periodic drive

Jordan-Wigner transformation

$$\hat{H}(t) = \sum_j (-J \hat{\sigma}_j^z \hat{\sigma}_{j+1}^z + h(t) \hat{\sigma}_j^x) \quad \Rightarrow \quad H = \sum_j J (c_i + c_i^\dagger)(c_{i+1} - c_{i+1}^\dagger) + h(t) c_i^\dagger c_i$$



End of the story?

Khemani, Lazarides, Moessner, Sondhi PRL 2015

Else, Bauer, Nayak PRL 2016

The problem : heating

Break integrability = Add interactions



Thermalize (Eigenstate thermalization hypothesis)



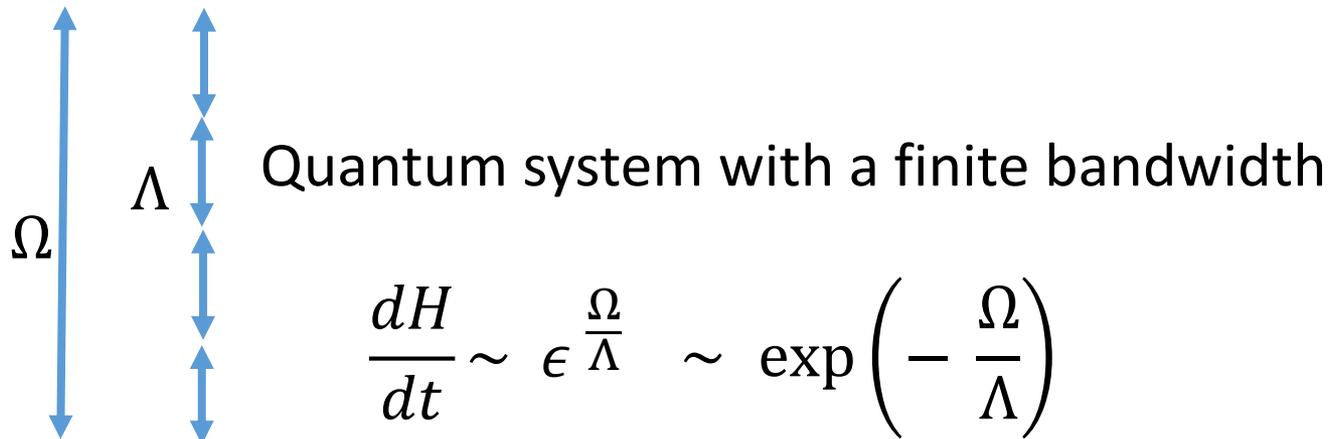
Steady state = infinite temperature

$$Z = e^{\beta H_F} \rightarrow e^{0 H_F} = 1$$



How to prevent heating?

- ✓ Integrability
- ✓ Disorder (many-body localization)
- ✓ **High-frequency drive (prethermalization)**



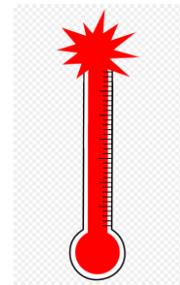
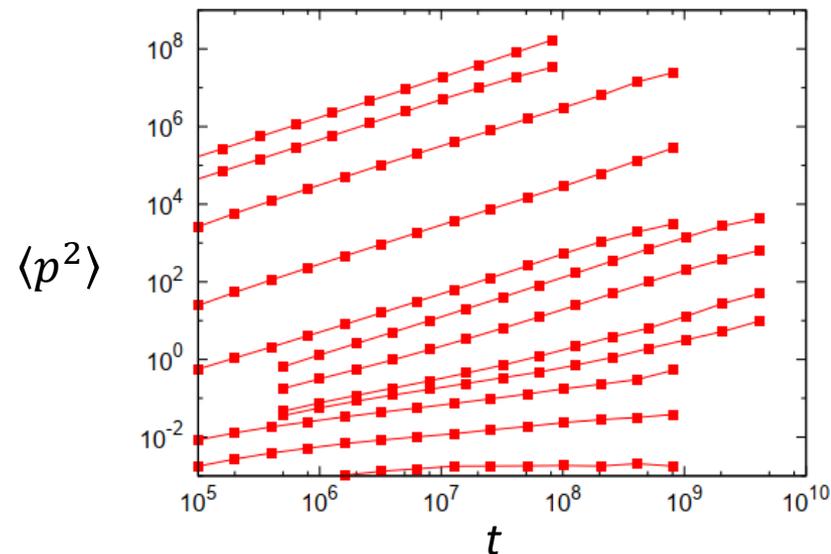
Quantum system with a finite bandwidth

$$\frac{dH}{dt} \sim \epsilon \frac{\Omega}{\Lambda} \sim \exp\left(-\frac{\Omega}{\Lambda}\right)$$

Abanin *et al* (PRL 2015) , (CMP 2017) - Kuwahara, Mori, Saito, (Ann.of. Phys 2015), (PRL 2015)

Classical systems with infinite bandwidth?

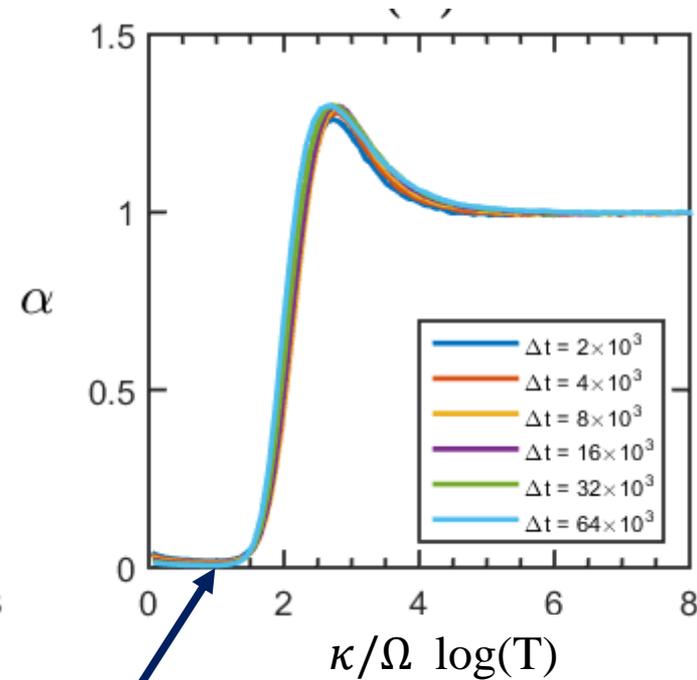
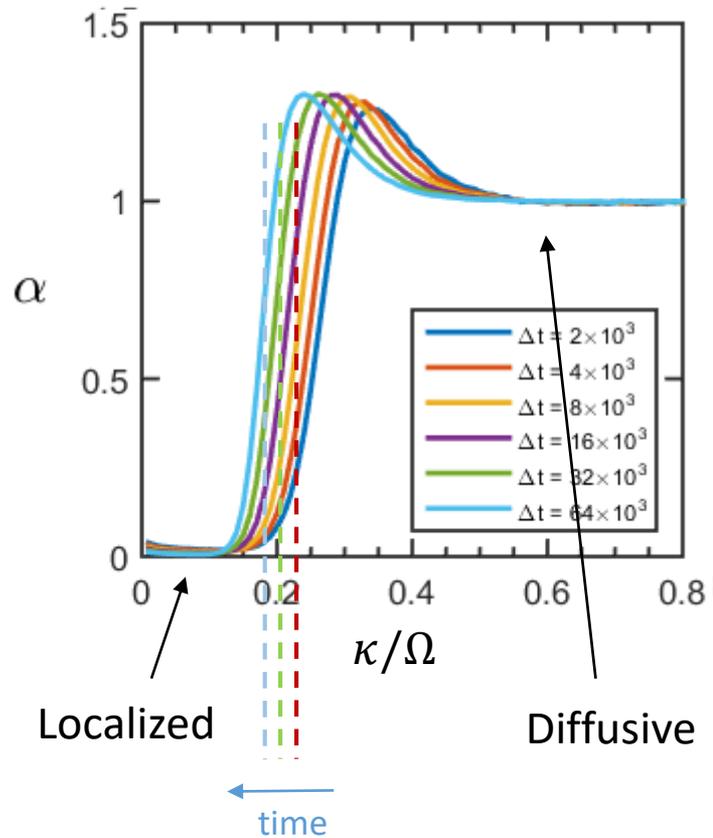
$$H = \sum_{j=1}^N \left[\frac{p_j^2}{2} - \kappa \cos(\phi_j - \phi_{j+1}) \sum_{n=-\infty}^{+\infty} \delta(t - n\tau) \right].$$



Kaneko & Konishi (1989) - Chirikov & Vecheslavov (1997)- Mulanski, Ahnert, Pikovski, Shepelyansky (2011)

Coupled kicked rotors – our results

$$\langle p(T)^2 \rangle = A t^\alpha$$



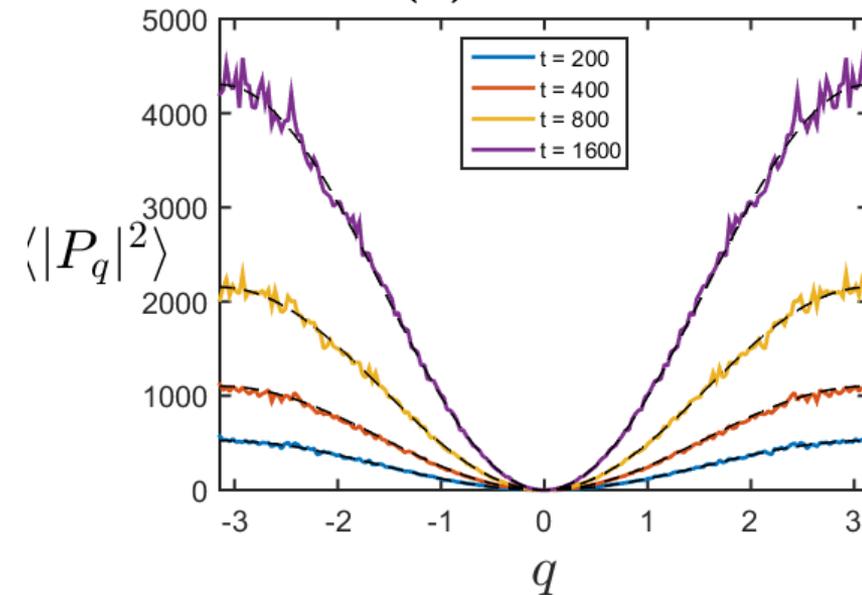
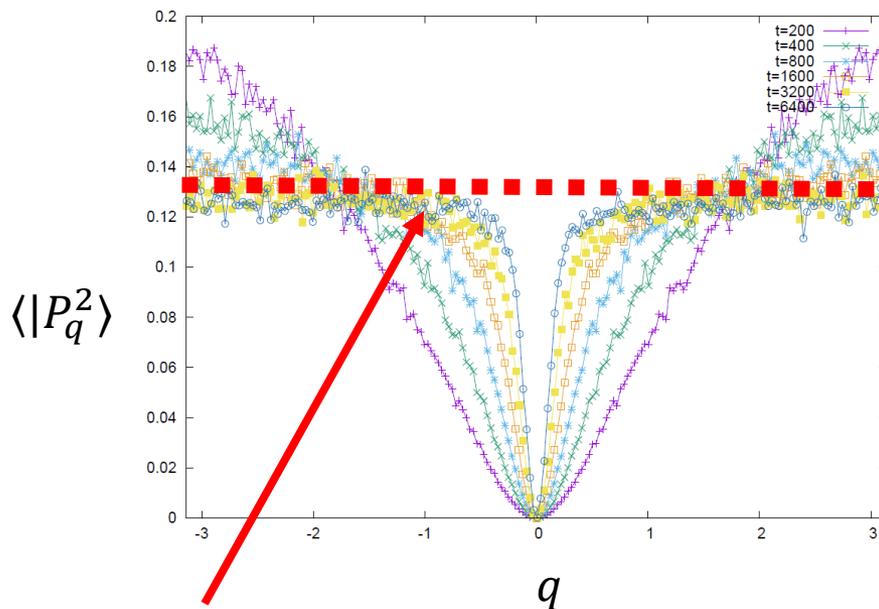
$$\frac{\kappa}{\Omega} \log(t) = 1.2$$

Coupled kicked rotors – momentum distrib.

$$P_q = \sum_n e^{iqn} \phi_n$$

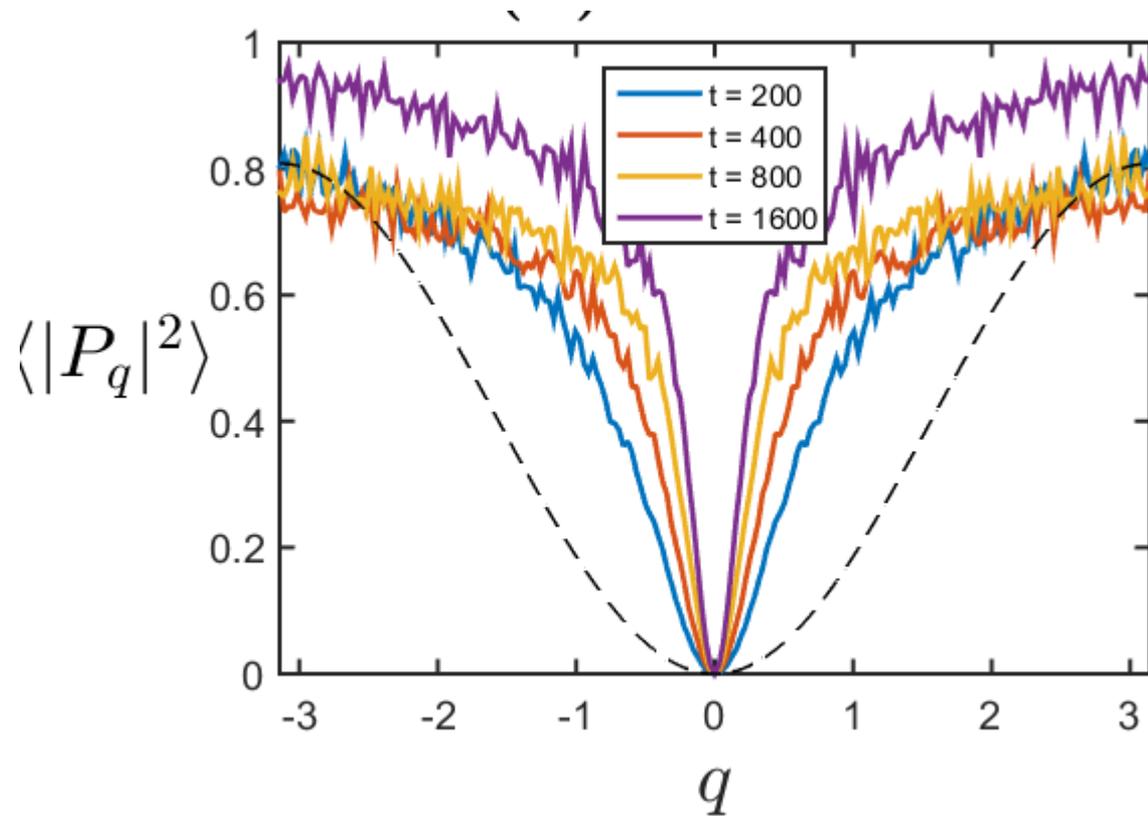
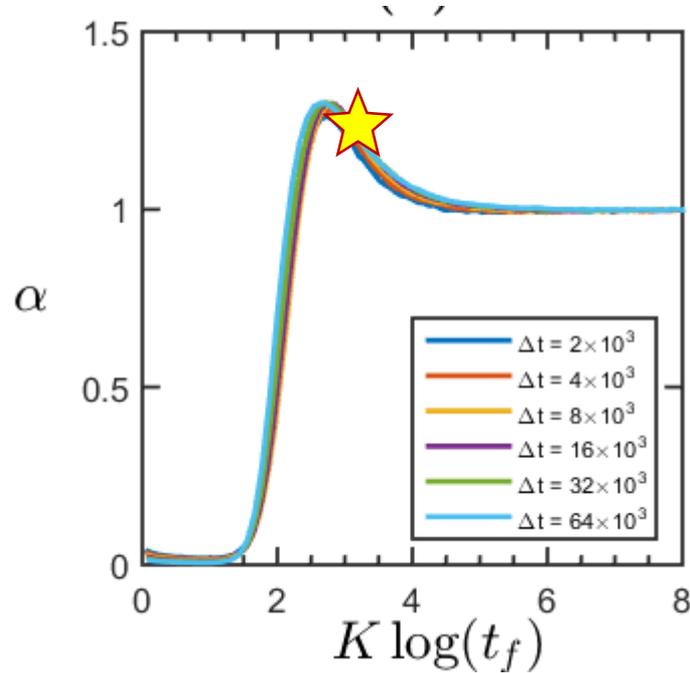
Localized

Diffusive



Prethermalization

Super-diffusion



Prethermalization in classical systems

$$H(t) = \sum_i \frac{p_i^2}{2} + K \Delta(t) \cos(\phi_i - \phi_{i+1})$$

$$\Delta(t) = \sum_n \delta(t - n \frac{2\pi}{\Omega})$$

Localizations threshold

$$\frac{\kappa}{\Omega} \log(\Omega t^*) = 1.2 \quad \longrightarrow \quad \Omega t^* = \exp\left(1.2 \frac{\Omega}{\kappa}\right)$$

$$\Omega \rightarrow \infty$$

~~High-frequency
suppresses heating
(Floquet Hamiltonian)~~

$$K \rightarrow 0$$

~~Reduced diffusion
close to integrability
(Arnold diffusion)~~

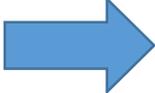
Prethermalization in classical systems : 1,2,3

1. Boltzmann distribution : $Z = \exp\left(-\frac{H_F}{T}\right)$

$$H_F = \sum_i \frac{p_i^2}{2} + \kappa\Omega \cos(\phi_i - \phi_{i+1})$$

2. Temperature = energy of the initial state : $T = \kappa\Omega$

3. Heating = resonance : $p_i - p_{i+1} = m \Omega$

 $\frac{dH}{dt} \sim \exp\left(-\frac{\Omega}{\kappa}\right)$

From Floquet Engineering to Pre-thermalization

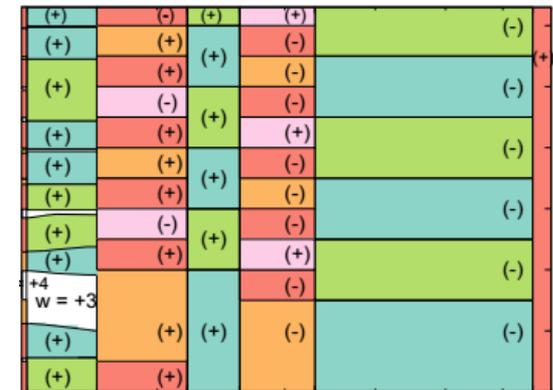
Periodically driven spin chain

- Periodically driven *quantum* Ising model

→ Time crystal & Floquet topological phases

Russomanno, Dalla Torre (EPL 2016)

Russomanno, Friedman, Dalla Torre (PRB 2017)



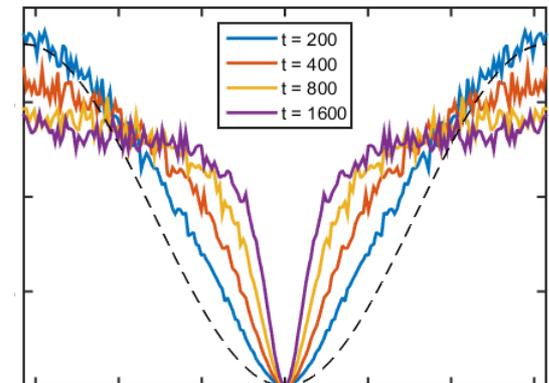
- Many body *classical* kicked rotor

→ Localization & Prethermalization

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Rajak, Citro, Dalla Torre (J. Phys. A: Math. Theor. 51 465001 (2018))

Rajak, Dana, Dalla Torre (in prep.)



Many-body quantum dynamics group



Classical Prethermalization

Dr. Atanu Rajak:
Dynamic localization of
a many-body kicked rotor



Quantum optics

Mor Roses:
Dicke superradiance
and counterlasing



Nonlinear optics (with Avi Pe'er)

Dr. Marcello Strinati:
Ising simulators with
parametric amplifiers



Topology in one dimension

Daniel Atzitz:
Symmetry resolved
entangled states



Superconducting circuits (with Michael Stern)

Inbar Shani:
Parametrically amplified
spin-cavity couplings



Strongly correlated materials

David Dentelski:
Short vs. Long-range
superconducting fluctuations

From Floquet Engineering to Pre-thermalization

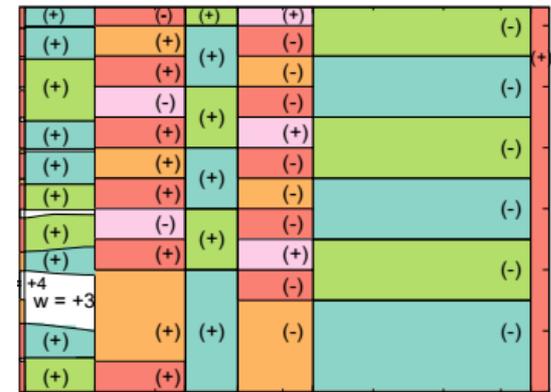
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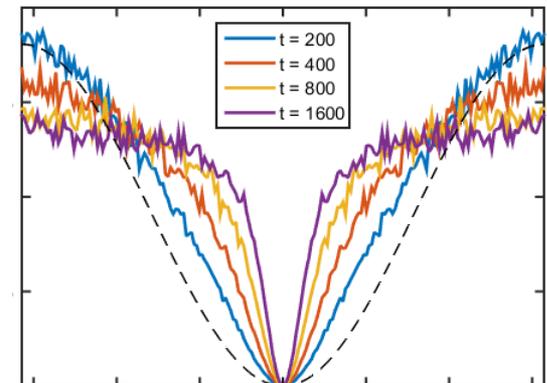
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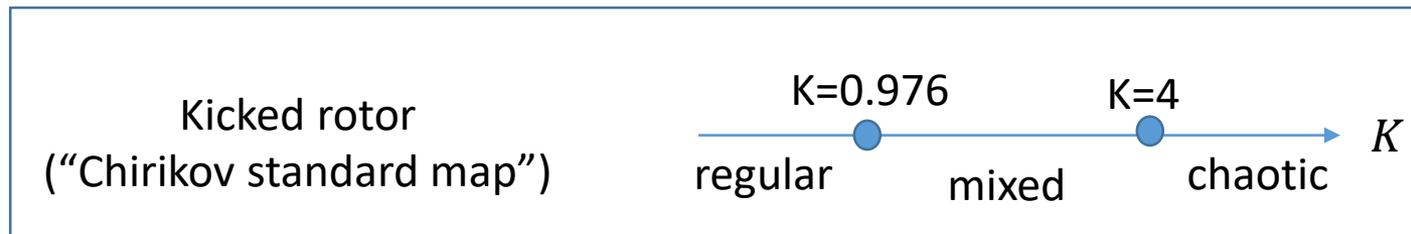


Extra slides

Coupled kicked rotors – quadratic expansion

$$H = \frac{1}{2} \sum_q \left[|P_q|^2 + K(q) |\phi_q|^2 \sum_{n=-\infty}^{+\infty} \delta(t - n\tau) \right],$$

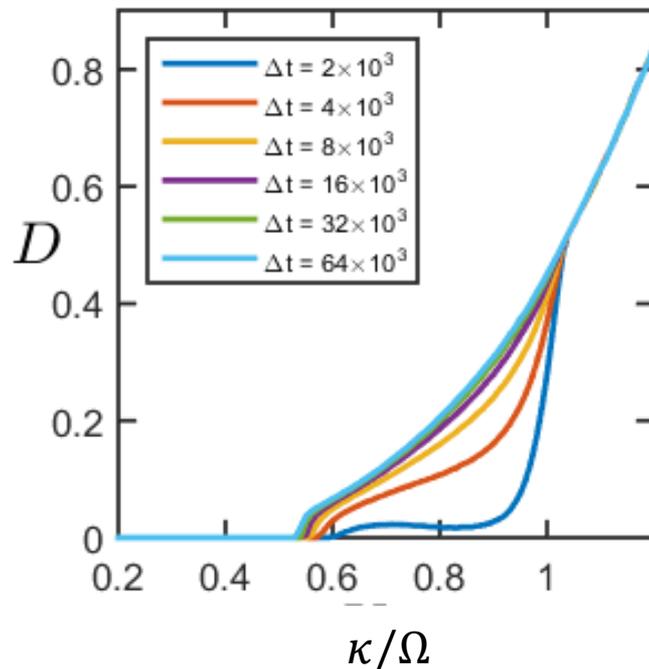
$$K(q) = \frac{4\kappa}{\Omega} \sin^2\left(\frac{q}{2}\right).$$



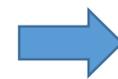
→ Transition at $\kappa/\Omega = 1$

Coupled kicked rotors – quadratic instability

Initial state: $\phi_j \approx 0$



$$H^4 = -\frac{\kappa}{24} \sum_{j=1}^N (\phi_j - \phi_{j+1})^4 \sum_{m=1}^{\infty} \cos\left(\frac{2\pi m}{\tau} t\right),$$

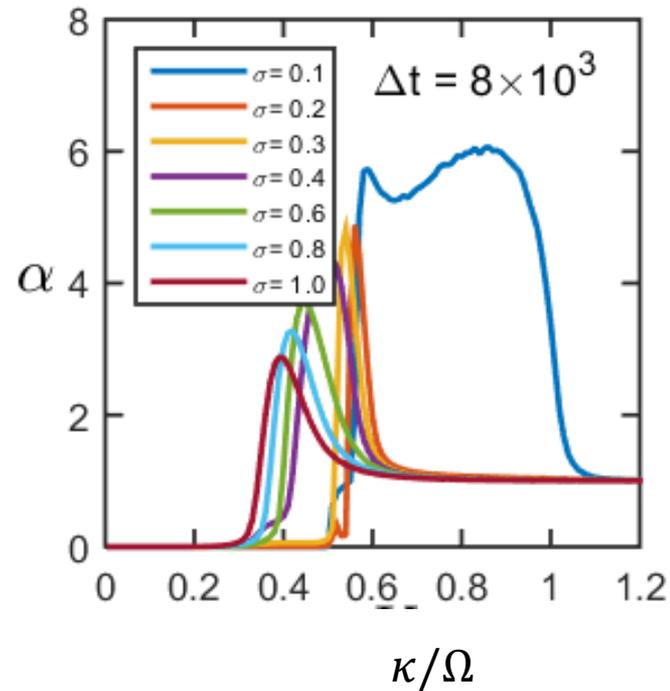
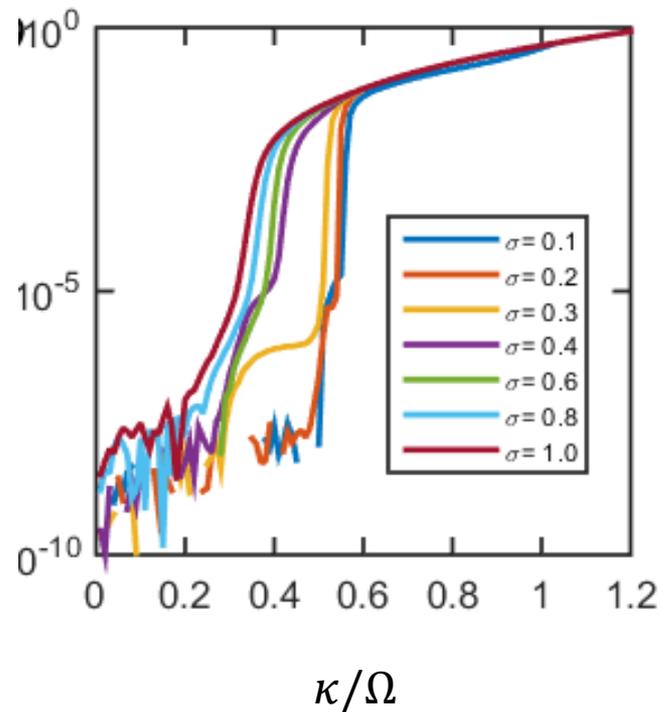


$$2\pi m \leq 4 \cos^{-1}(1 - 2\kappa/\Omega)$$

$$m = 4 \Rightarrow \frac{\kappa}{\Omega} > 0.5$$

Coupled kicked rotors – from quadratic to marginal

Initial state: $\langle \phi_j^2 \rangle = \sigma$



Finite size effect

