

What is the difference between Dicke superradiance and lasing?



P Kirton, M M Roses, J Keeling, and E G Dalla Torre,
“Introduction to the Dicke model: from equilibrium
to nonequilibrium, and vice versa.” *arXiv:1805.09828*

Nonequilibrium Quantum Dynamics



On this project:

Mor Roses

Marcello C. Strinati

Collaborators:

Eugene Demler

Mikhail Lukin

Yulia Schadilova

Hakan Tureci

Competitors:

Jonathan Keeling

Peter Kirton



<http://www.facebook.com/nonequilibrium>



What is the difference between Dicke superradiance and lasing?

A. SR is quantum and lasing is classical

B. Lasing is quantum and SR is classical

C. Only one of them is a phase transition

D. There is no real difference

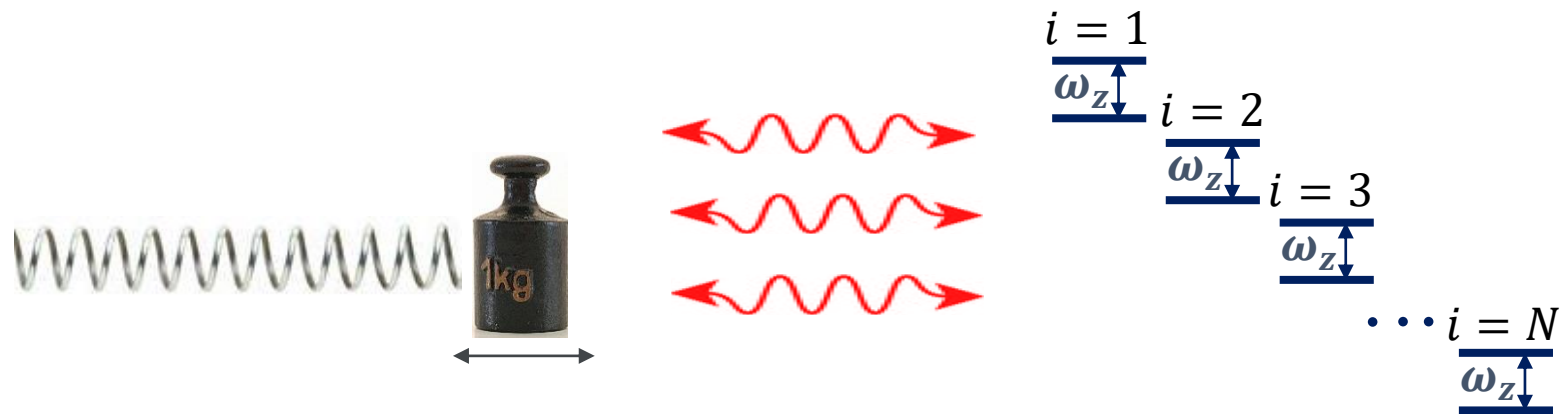
Outline

Dicke superradiance (8 slides)

Lasing (4 slides)

Extras (2 slides)

Dicke model



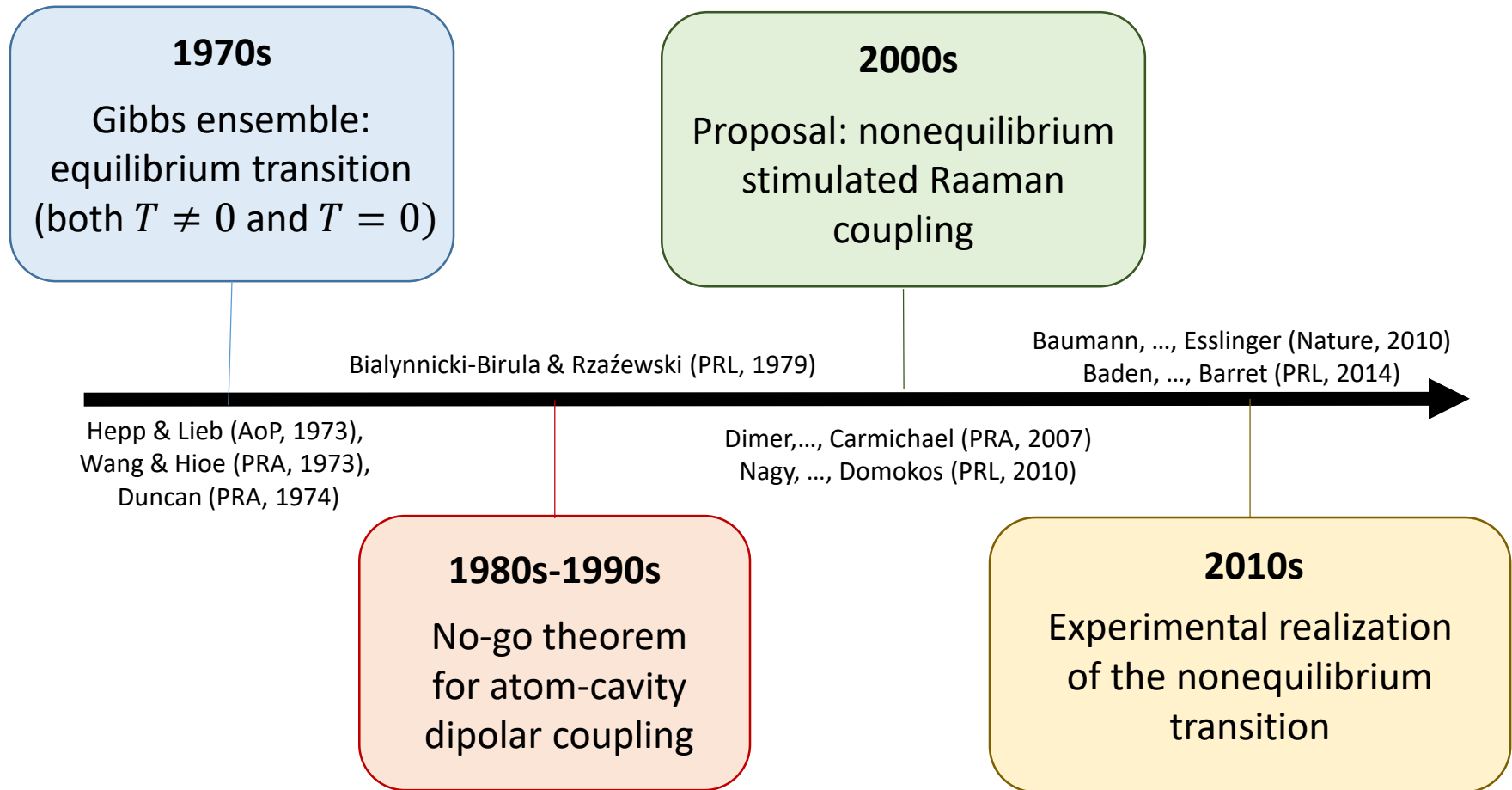
$$H = \frac{1}{2}(X^2 + P^2) + \frac{\lambda}{\sqrt{N}} \sum_{i=1}^N X \sigma_i^x + \omega_z \sum_{i=1}^N \sigma_i^z$$

“Ising model of quantum optics”

Credit: MASS - <http://www.fotocommunity.de/photo/1-kilogramm-kajo-schmitz/20034234>

SPRING <http://depositphotos.com/10120691/stock-photo-metal-spring-isolated-on-white.html>

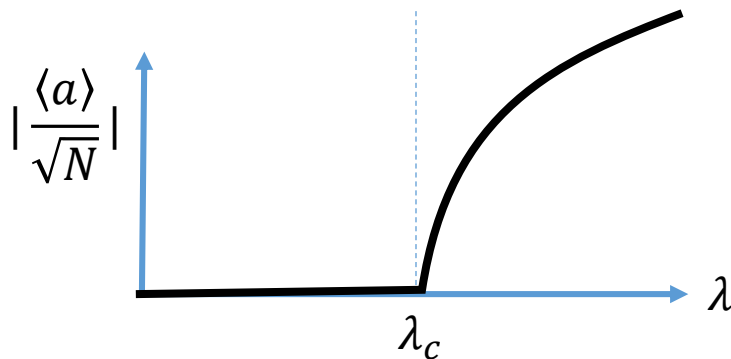
Timeline of Dicke superradiance



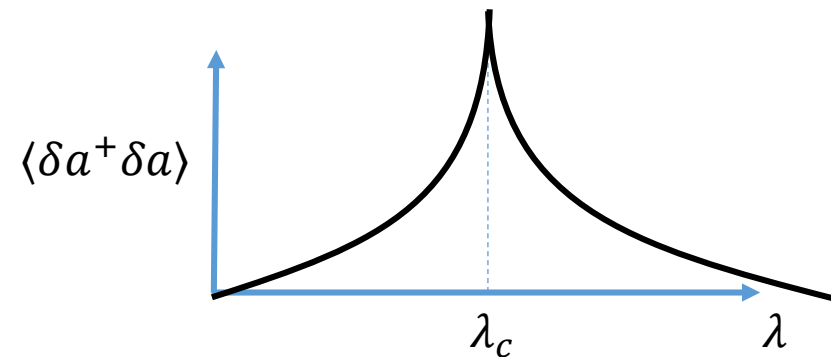
Dicke superradiant transition

$$H = \omega_0 a^\dagger a + \frac{\lambda}{\sqrt{N}} (a + a^\dagger) S^x + \omega_z S^z$$

Conservation of the total spin $\vec{S} = \sum_i \vec{\sigma}_i$



Ising mean field
 $\langle a \rangle = |\lambda - \lambda_c|^{-1/2}$



“Flux exponent”
 $\langle \delta a^\dagger \delta a \rangle = |\lambda - \lambda_c|^{-1}$

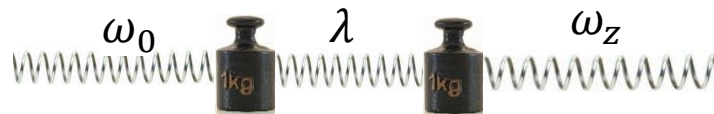
Tureci & Domokos groups (2014)

Dicke superradiance – critical theory

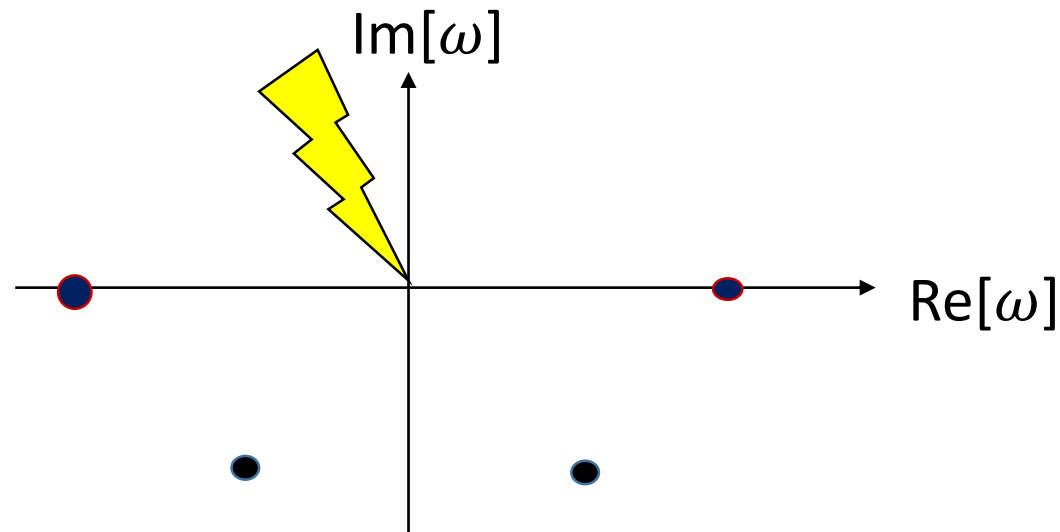
$$H = \omega_0 a^\dagger a + \frac{\lambda}{\sqrt{N}} (a + a^\dagger) S^x + \omega_z S^z$$

Holstein-Primakoff (linearization)

$$S^+ \approx \sqrt{N} b^+$$



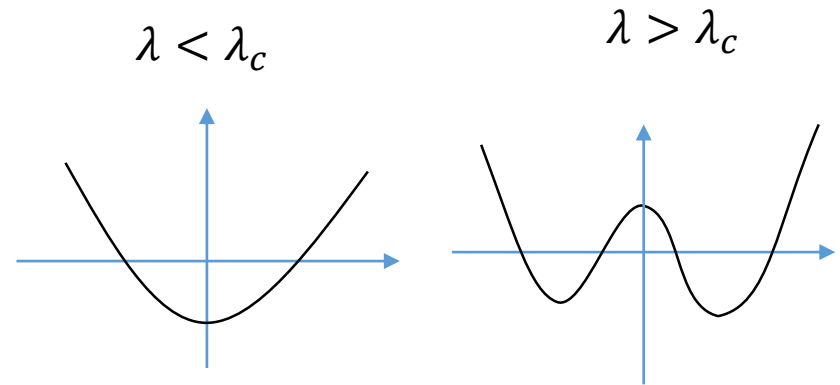
$\lambda = 0$



Dicke superradiance – critical theory

Mean field phase transition (Landau)

$$H = \frac{p^2}{2m} + \frac{1}{2}(\lambda - \lambda_c)x^2 + \beta x^4$$



Critical theory:

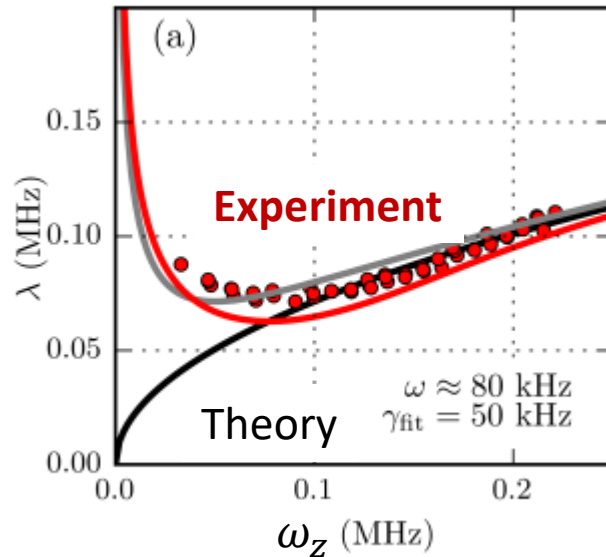
$$H = \frac{p^2}{2m} + \frac{1}{2}(\lambda - \lambda_c)x^2$$

Equipartition: $\langle (\lambda - \lambda_c)x^2 \rangle = k_B T \Rightarrow \langle x^2 \rangle = \frac{k_B T}{|\lambda - \lambda_c|}$

Dalla Torre et al PRA (2013)

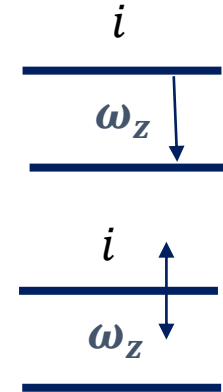
Dicke transition - revisited

Zhiqiang et al, PRA (2018)



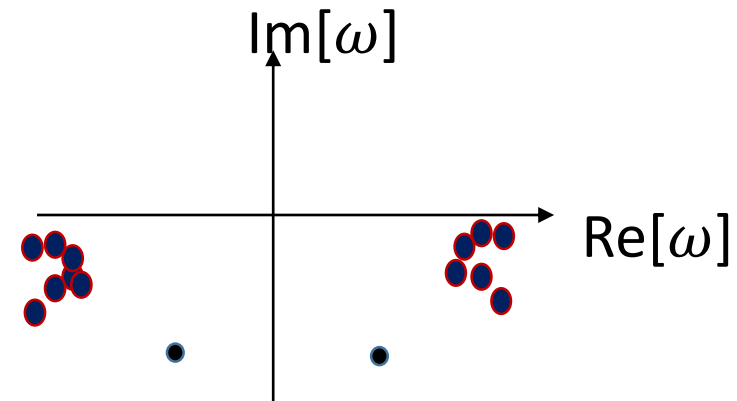
Single-atom decay

Single-atom dephasing
(transverse Doppler shift)



Total spin not conserved

$$H = \omega_0 a^\dagger a + \frac{\lambda}{\sqrt{N}} (a + a^\dagger) \sum_{i=1}^N \sigma_i^x + \omega_z \sum_{i=1}^N \sigma_i^z$$



Diagrammatic approach to the Dicke transition

Fermionic representation of spin-1/2 $\sigma_i^z = f_i^\dagger f_i - \frac{1}{2}$, $\sigma_i^- = \eta_i f_i$

Bare terms:



Leading 1/N term:



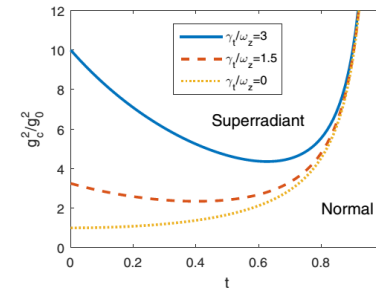
*** controlled 1/N expansion ***

Result 1 – critical coupling

$$\lambda_c^{-2} = \frac{2}{N} \sum_{i=1}^N \int_0^\infty dt \langle [\sigma_i^x(0), \sigma_i^x(t)] \rangle$$

- ✓ Equilibrium ($T = 0$ and $T \neq 0$)
- ✓ Nonequilibrium with spin conservation

NEW: Nonequilibrium without spin conservation



Polarization = sufficient condition for Dicke transition
(no correlations or entanglement)

- Cfr. Lamb theory of lasing transition

Dalla Torre et al PRA (2017), Kirton and Keeling PRL (2017)

Outline

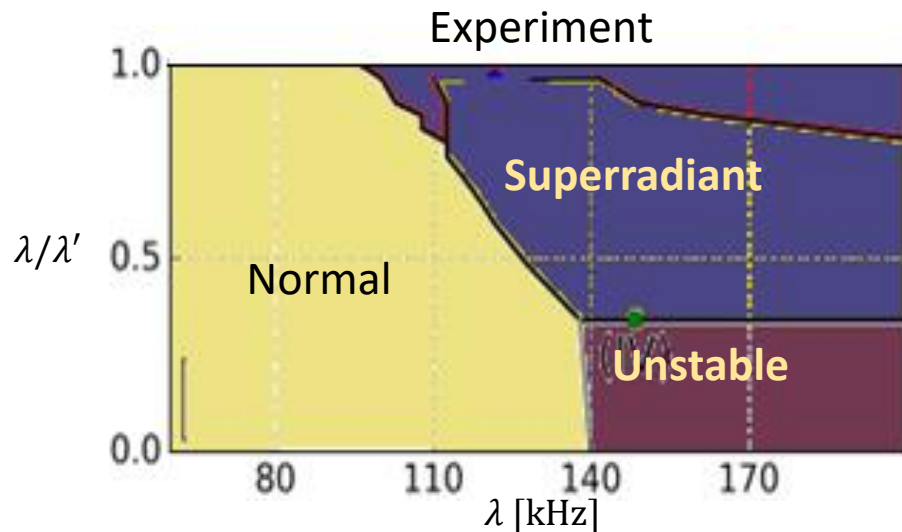
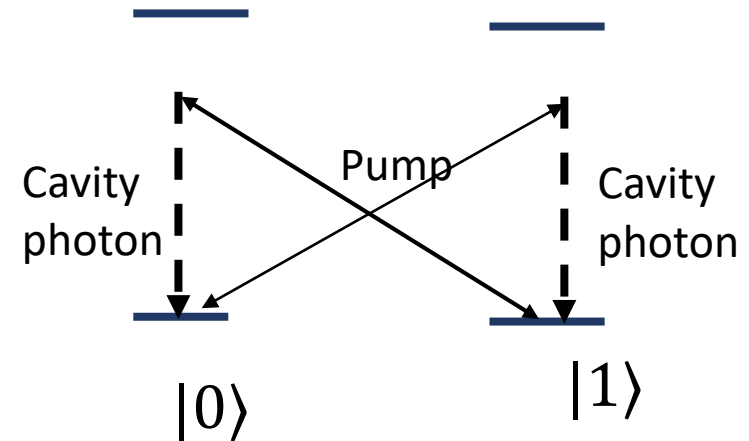
Dicke superradiance (8 slides)

Lasing (4 slides)

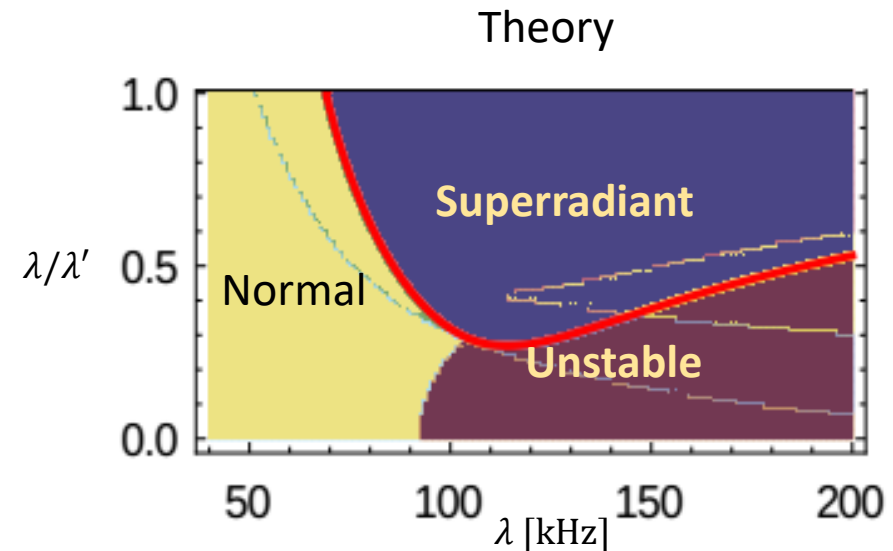
Extras (4 slides)

Generalized Dicke model

$$H = \omega_0 a^\dagger a + \omega_z \sum_{i=1}^N \sigma_i^z + \frac{\lambda}{\sqrt{N}} a \sum_{i=1}^N \sigma_i^+ + \frac{\lambda'}{\sqrt{N}} a^\dagger \sum_{i=1}^N \sigma_i^+ + H.c.$$



Zhiquiang et al Optica (2017)

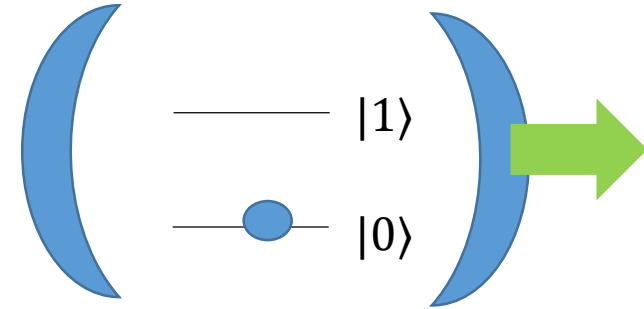


Sh

Presented in Cargese on May 11, 2017.
cfr. With Kirton, Keeling, arxiv:1710.06212

Anti-Jaynes-Cummings model

$$H = \omega_0 a^\dagger a + \frac{\lambda'}{\sqrt{N}} \sum_{i=1}^N a^\dagger \sigma_i^+ + a \sigma_i^- + \omega_z \sum_{i=1}^N \sigma_i^z$$



Threshold: Fermi Golden rule

$$\frac{(\lambda')^2}{\omega_0 - \omega_z} = \kappa$$

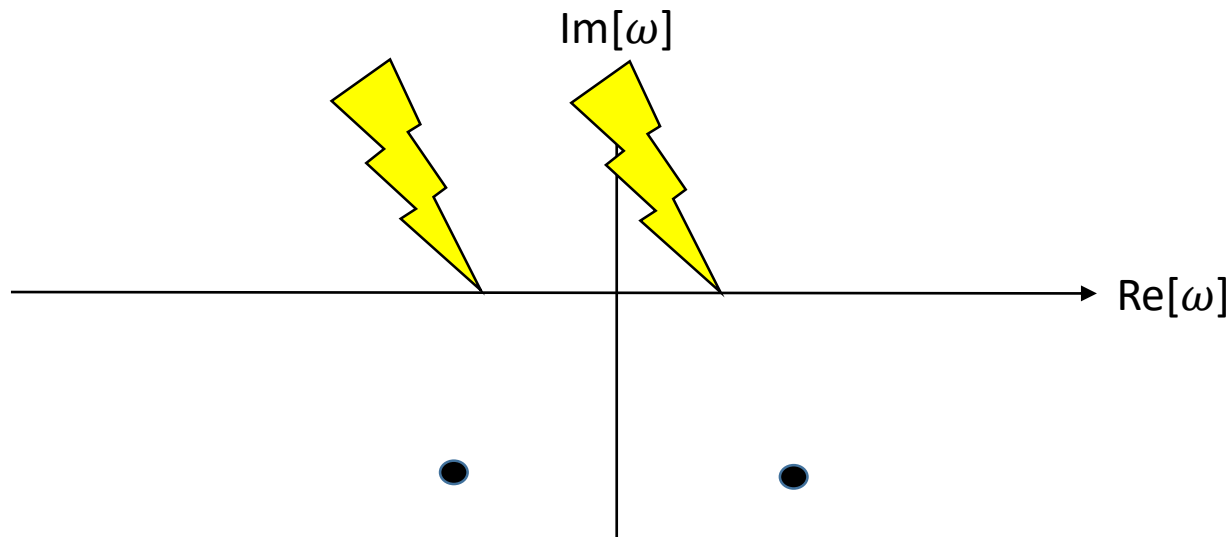
Decay $|1\rangle \rightarrow |0\rangle$ = “counter-repumping”

→ Counterlasing

Kirton, Keeling, NJP (2018), Shchadilova et al arxiv (2018)

Counterlasing transition – critical theory

$$H = \omega_0 a^\dagger a + \frac{\lambda'}{\sqrt{N}} a^\dagger \sum_{i=1}^N \sigma_i^+ + \omega_z \sum_{i=1}^N \sigma_i^z$$

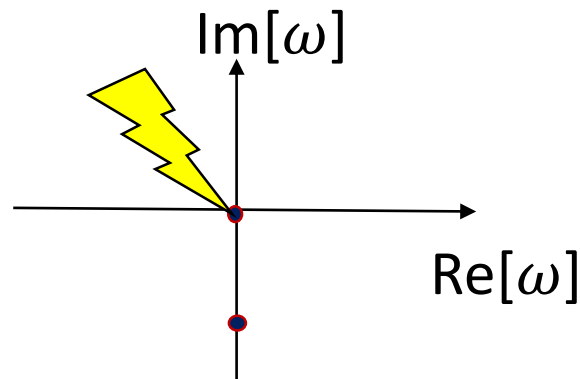


- No Diverging fluctuations
- Phase degree of freedom



Dicke superradiance vs Lasing

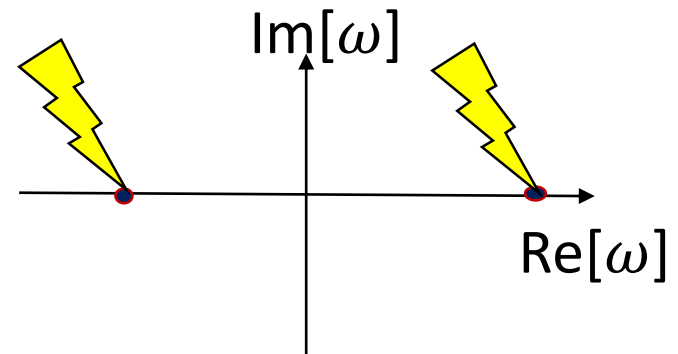
Pitchfork



- Divergence
- Stationary

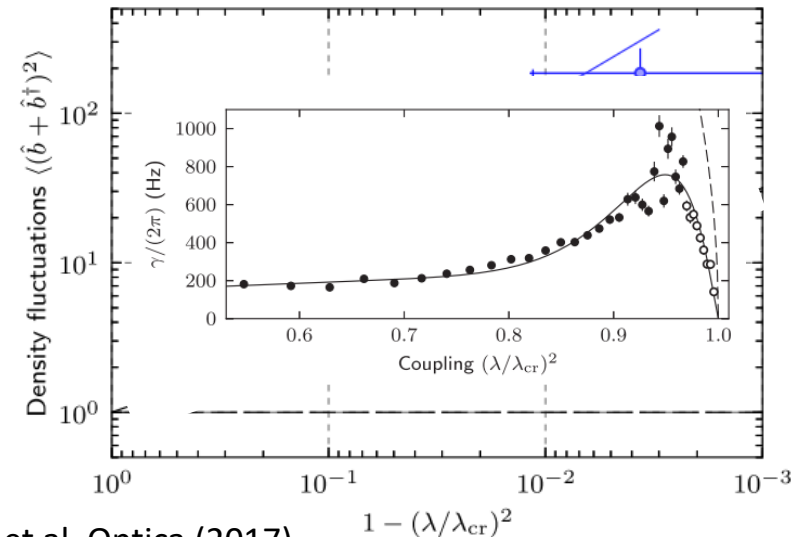
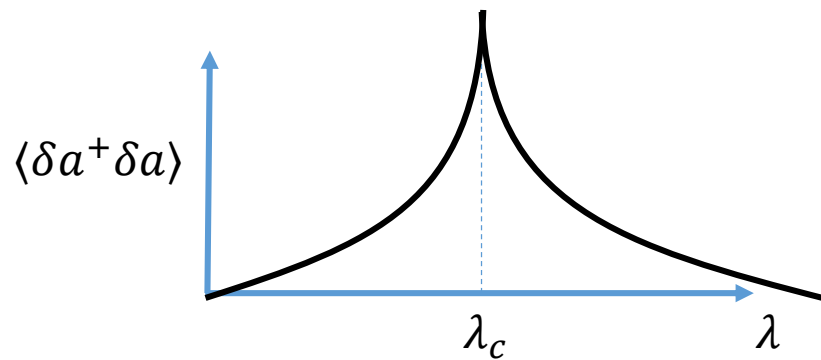
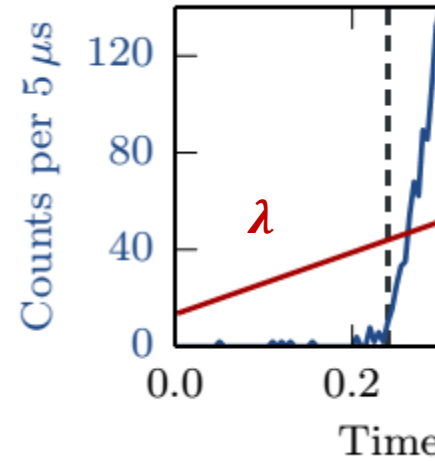
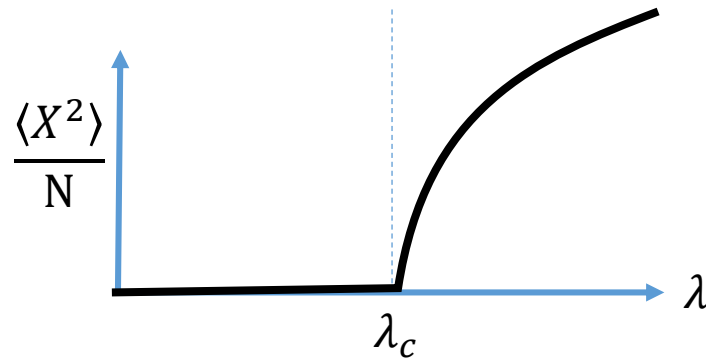


Hopf instability



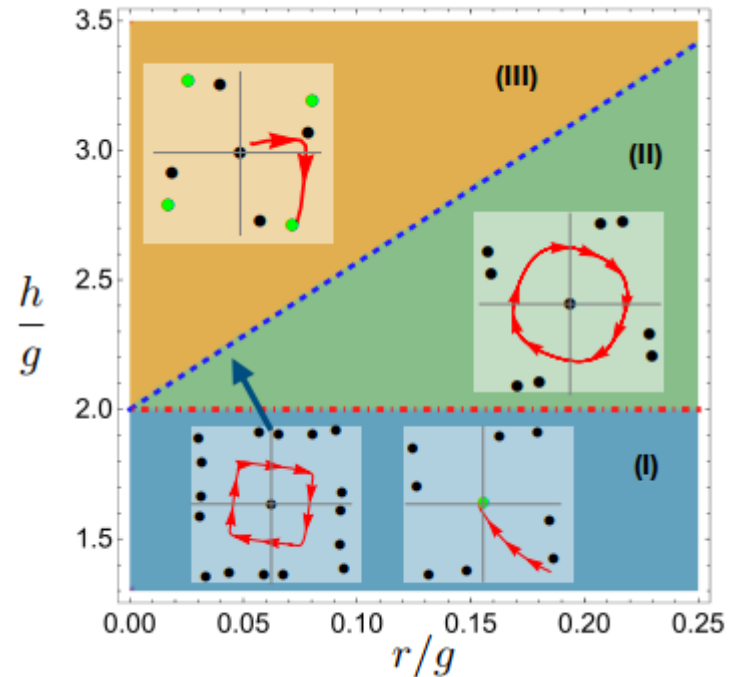
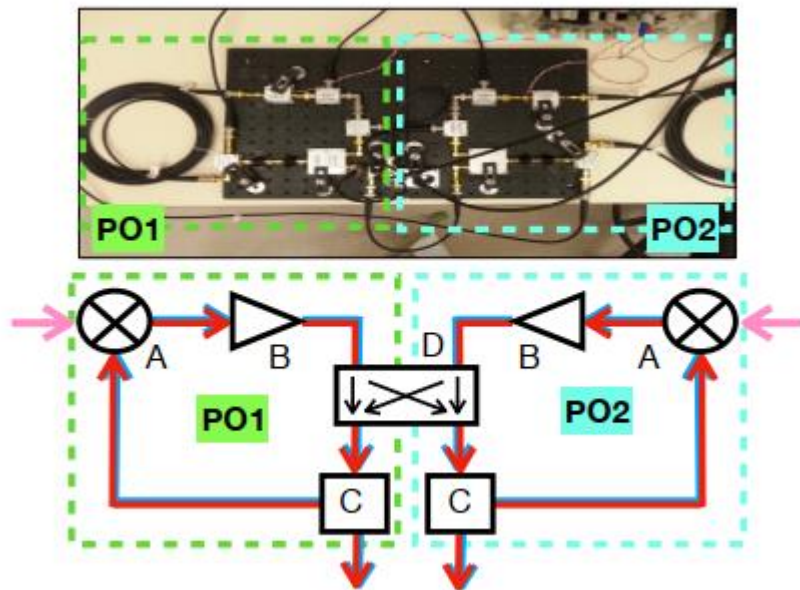
me)

The end of the story?



Brennecke et al PNAS (2013), Zhiqiang et al, Optica (2017)

Another system with the same physics



L. Bello, M. Calvanese Strinati, E. G. Dalla Torre, A. Pe'er (in preparation)

Non-equilibrium Quantum Dynamics



Periodically-driven systems

Dr. Atanu Rajak:
Dynamic localization of
a many-body kicked rotor



Topological phases

Ben Yamin:
Measuring entanglement
in quantum computers

Dynamics of closed systems



Meitar Goldferb:
Multicomponent
ultracold atoms



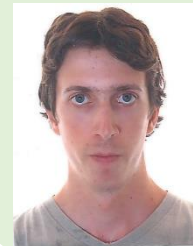
Quantum optics

Mor Roses:
Real-time dynamics
of the extended Dicke model



Strongly correlated materials

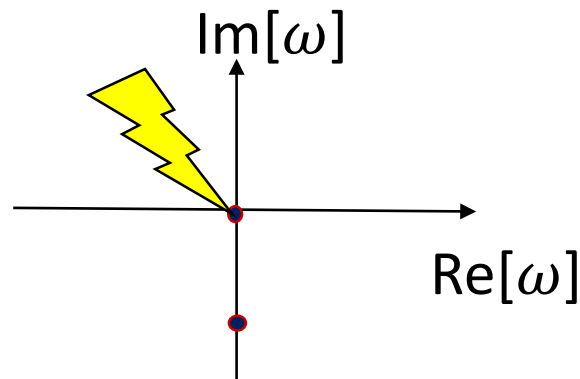
David Dentelski:
Short vs. Long-range
superconducting fluctuations



Dr. Marcello Strinati:
Simulation of spin systems
in (quantum) optics

Dicke superradiance vs Lasing

Pitchfork



Hopf instability

